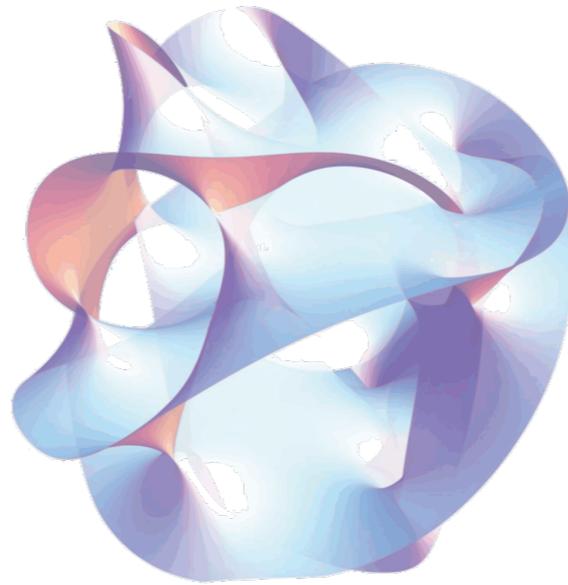
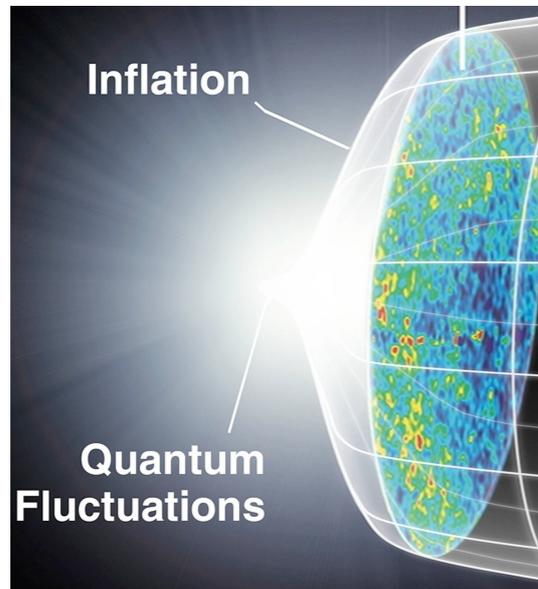


# Conformal Bootstrap and Quantum Field Theory

Modern Trends in Particle Physics  
NathFest 2019

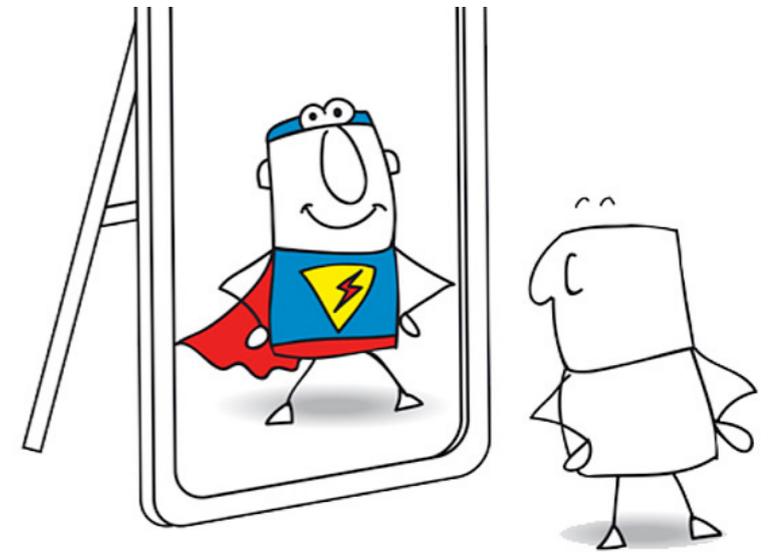
# Modern Trends in Particle Physics

inflation

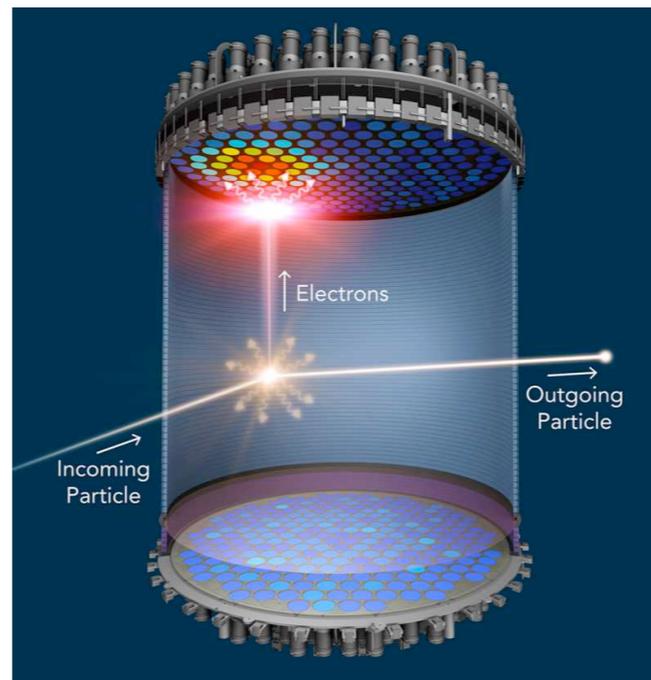
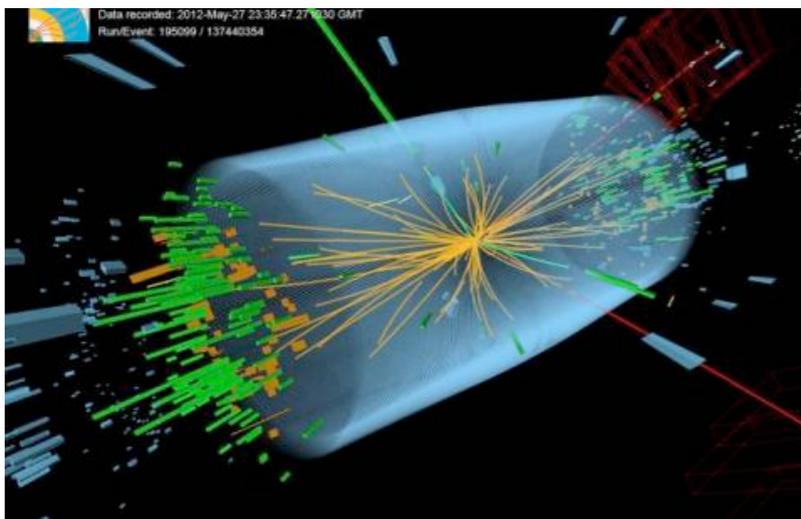


dark matter

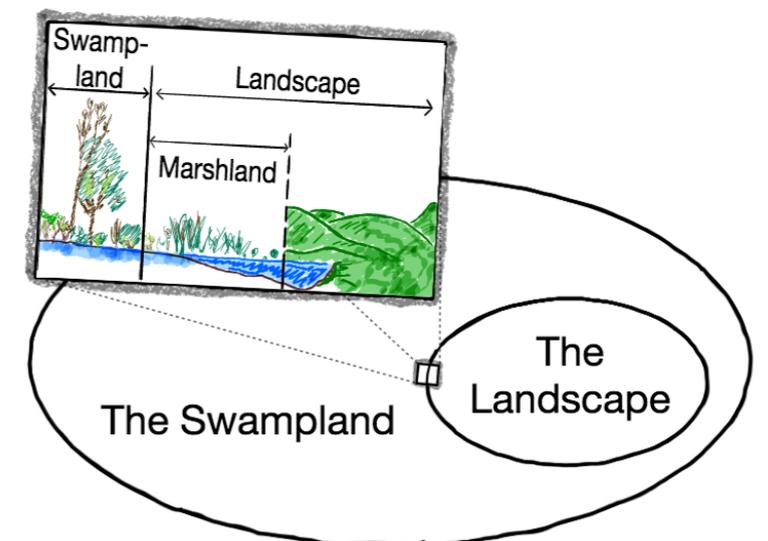
supersymmetry



colliders

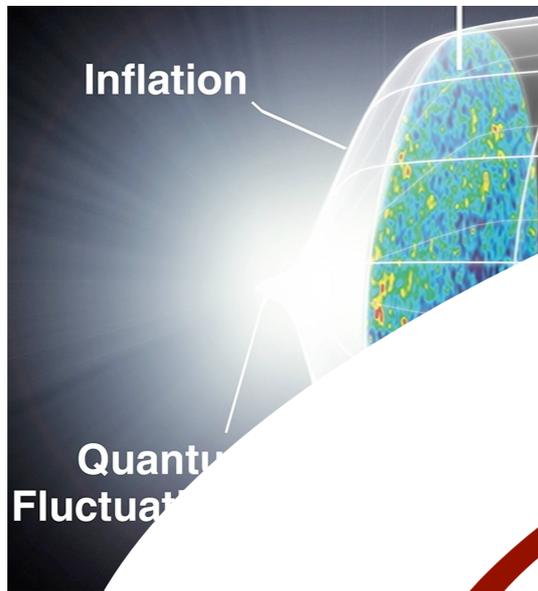


landscape



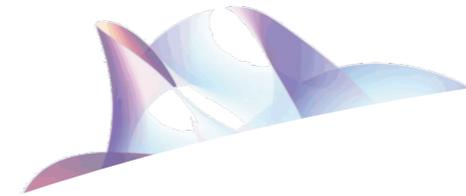
# Modern Trends in Particle Physics

inflation



supersymmetry

... |||

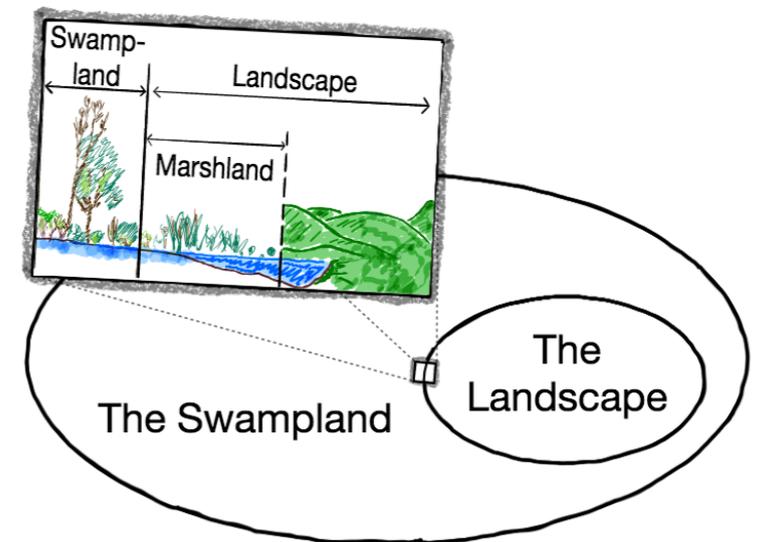
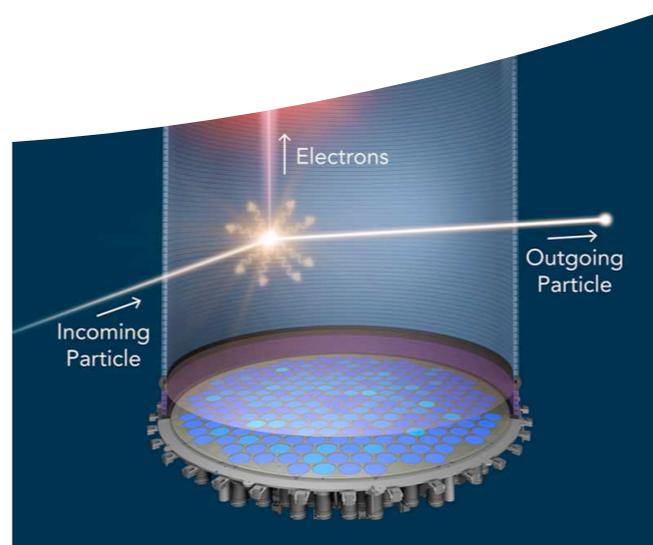
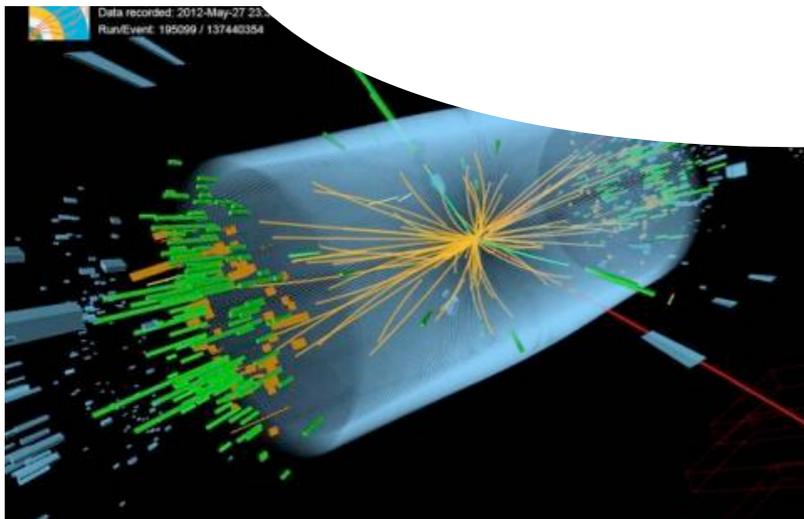


It's all just

QFT



landscape



It's all just

QFT = CFT + EFT

*Philosophy Master:* Because there is no other way to express oneself but through prose or poetry. Whatever is not prose, is poetry, and whatever is not poetry is prose.

*Jourdain:* When I say, Nicole! Bring me my slippers, is that prose?

*Philosophy Master:* Yes, sir.

*Jourdain:* So I have been speaking prose for years without even knowing it!



It's all just

QFT = CFT + EFT

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What is  
Quantum  
Field Theory?

# What is Quantum Field Theory?

There are a number of different  
nonperturbative answers

- 1) Lattice Path Integral
- 2) S-Matrix
- 3) Conformal Bootstrap

# What is Quantum Field Theory?

## 1) Lattice Path Integral

QFT is given by Euclidean path integral

Advantages

Disadvantages

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## 1) Lattice Path Integral

QFT is given by Euclidean path integral

### Advantages

Well-developed tool, gives concrete numeric results for many physical quantities at strong coupling

Gives a rigorous definition of many QFTs

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## 1) Lattice Path Integral

QFT is given by Euclidean path integral

### Advantages

Well-developed tool, gives concrete numeric results for many physical quantities at strong coupling

Gives a rigorous definition of many QFTs

### Disadvantages

Difficult/impossible to get many Lorentzian quantities

Latticization not always possible (e.g. non-Lagrangian theories, or chiral gauge theories, esp. **Standard Model!**)

The real world is not a lattice

# What is Quantum Field Theory?

## 2) S-Matrix

QFT is given by S-matrix satisfying various consistency conditions (crossing, “analyticity”, unitarity)

Advantages

Disadvantages

# What is Quantum Field Theory?

## 2) S-Matrix

QFT is given by S-matrix satisfying various consistency conditions (crossing, “analyticity”, unitarity)

### Advantages

Focuses on physically relevant quantities - useful & conceptually satisfying

Many beautiful and useful patterns emerge

“Explains” why QFT is necessary (Weinberg)

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### Advantages

Focuses on physically relevant quantities - useful & conceptually satisfying

Many beautiful and useful patterns emerge

“Explains” why QFT is necessary (Weinberg)

### Disadvantages

Difficult/impossible to construct exact S-matrix beyond perturbation theory

Not all theories have particles and S-matrices

More to life than scattering

# What is Quantum Field Theory?

## 3) Conformal Bootstrap

CFTs=solutions to Bootstrap Equation

QFT is given by RG flows between CFTs

Advantages

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CFTs are nonperturbatively defined in terms of CFT “data” and bootstrap equation

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Gives Lorentzian and Euclidean correlators

No need of a Lagrangian or particles!

### Disadvantages

Not always possible to find CFT data in practice

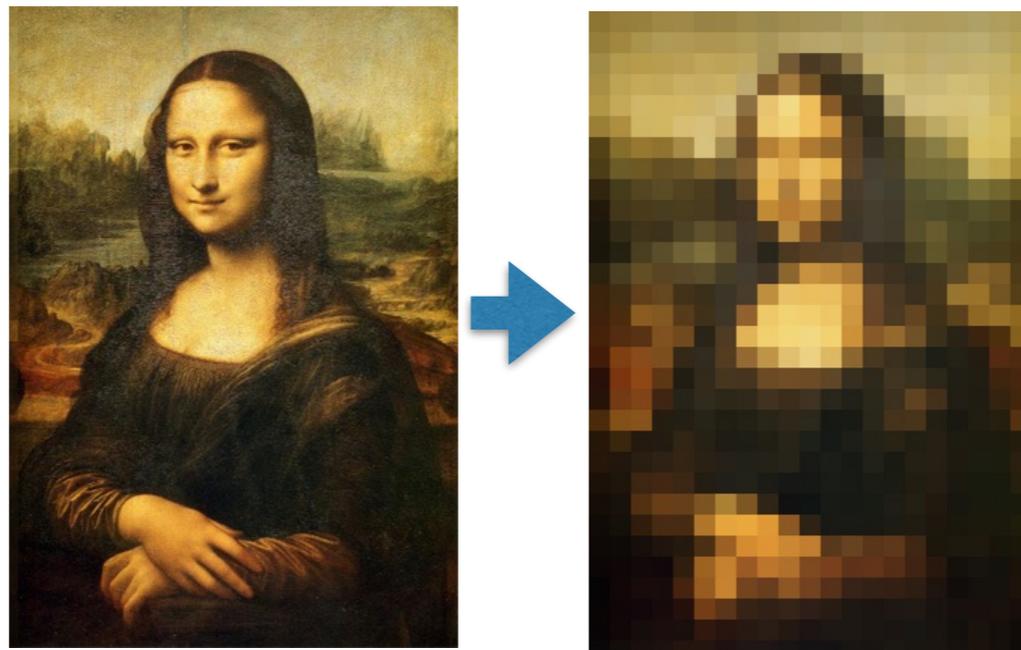
Not all QFTs are CFTs in UV  
Sometimes all we have is EFT  
e.g. gravity, string theory\*

\*But string theory = CFT on world sheet,  
and gravity in AdS=CFT,  
so still just CFT in the UV?

# RG Flow

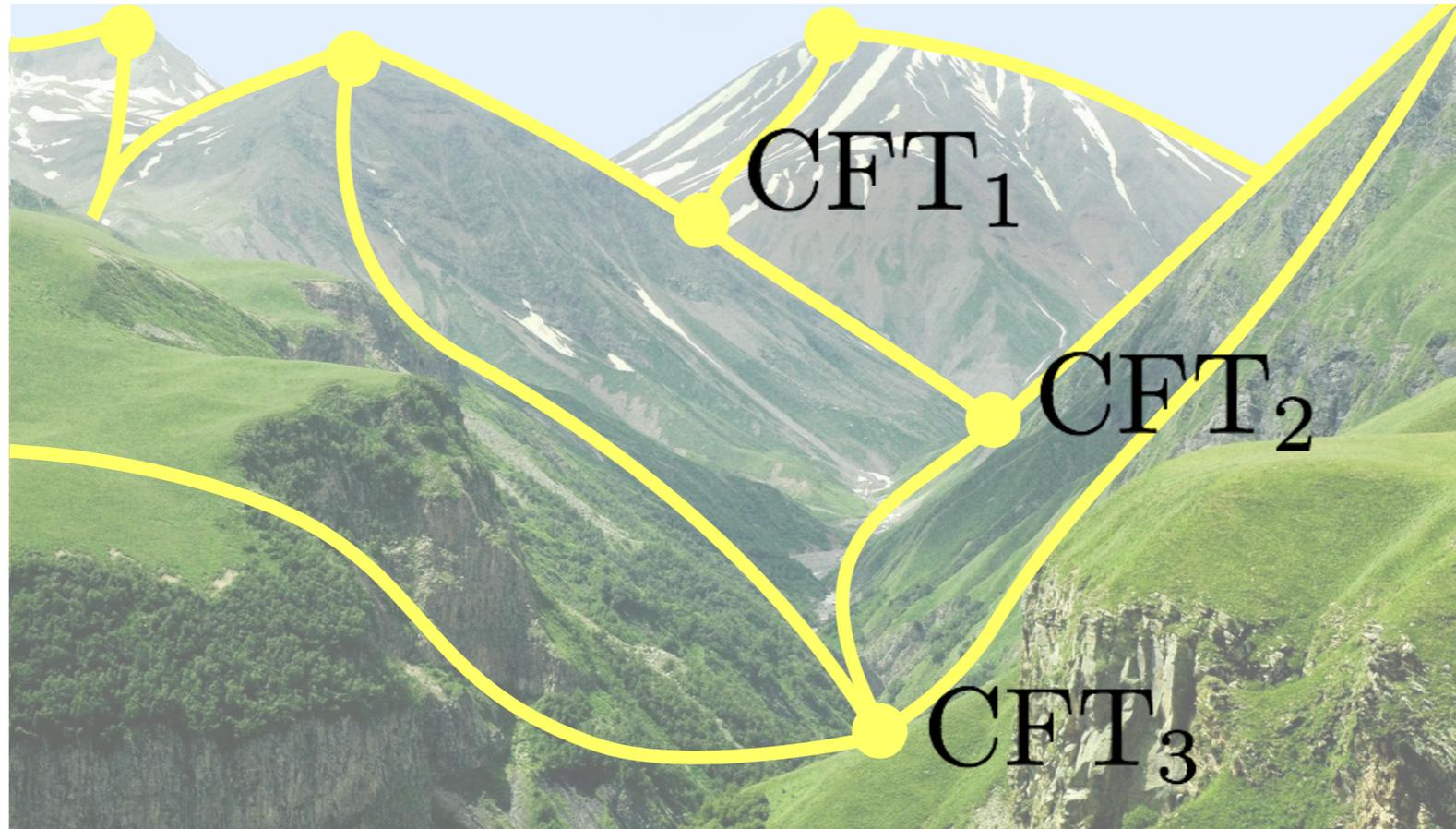
Ken Wilson :

Organize by energy scale  $=(\text{distance scale})^{-1}$

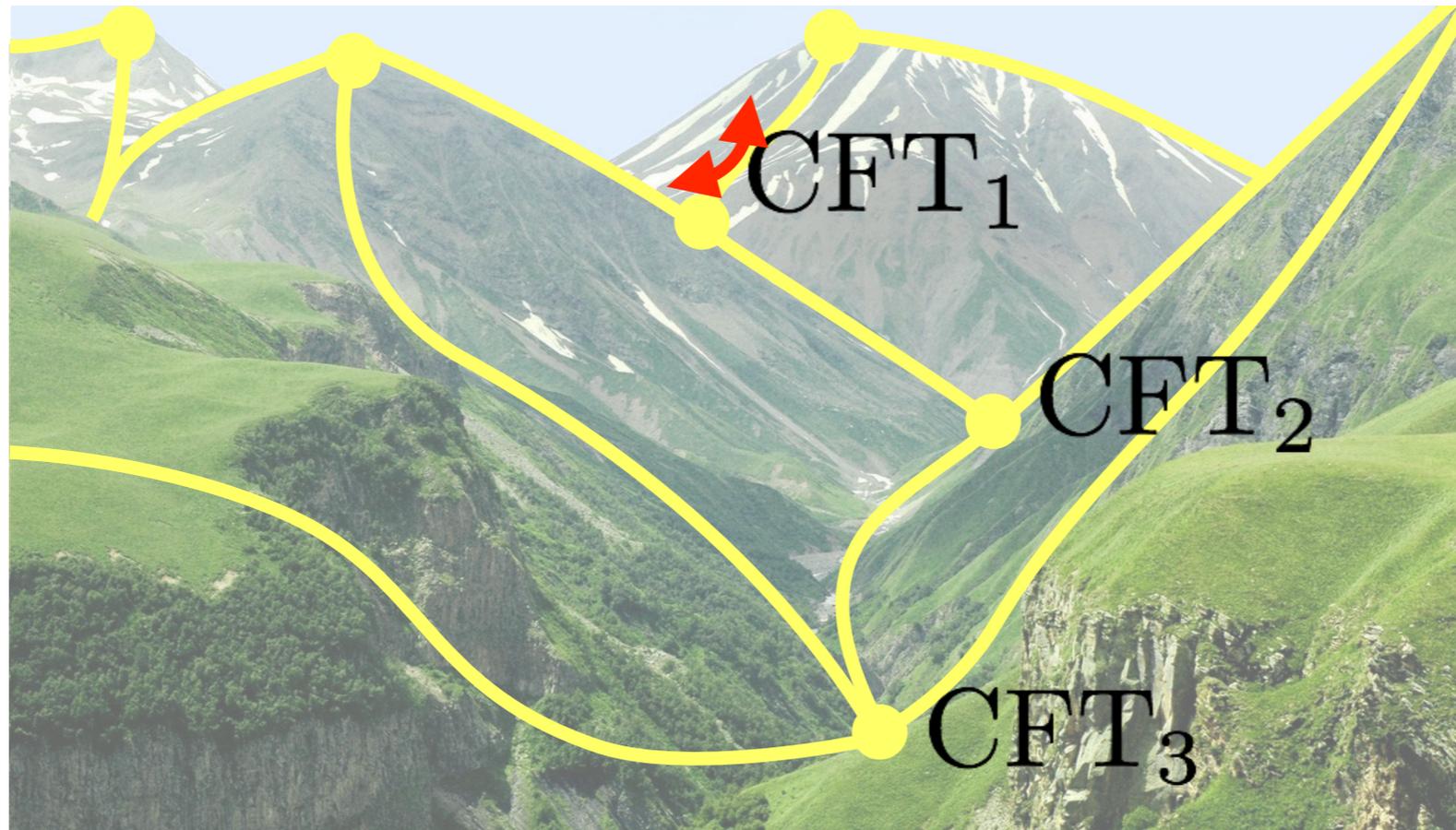


Classify terms in Hamiltonian as “relevant” or “irrelevant”

# RG Flow between CFTs

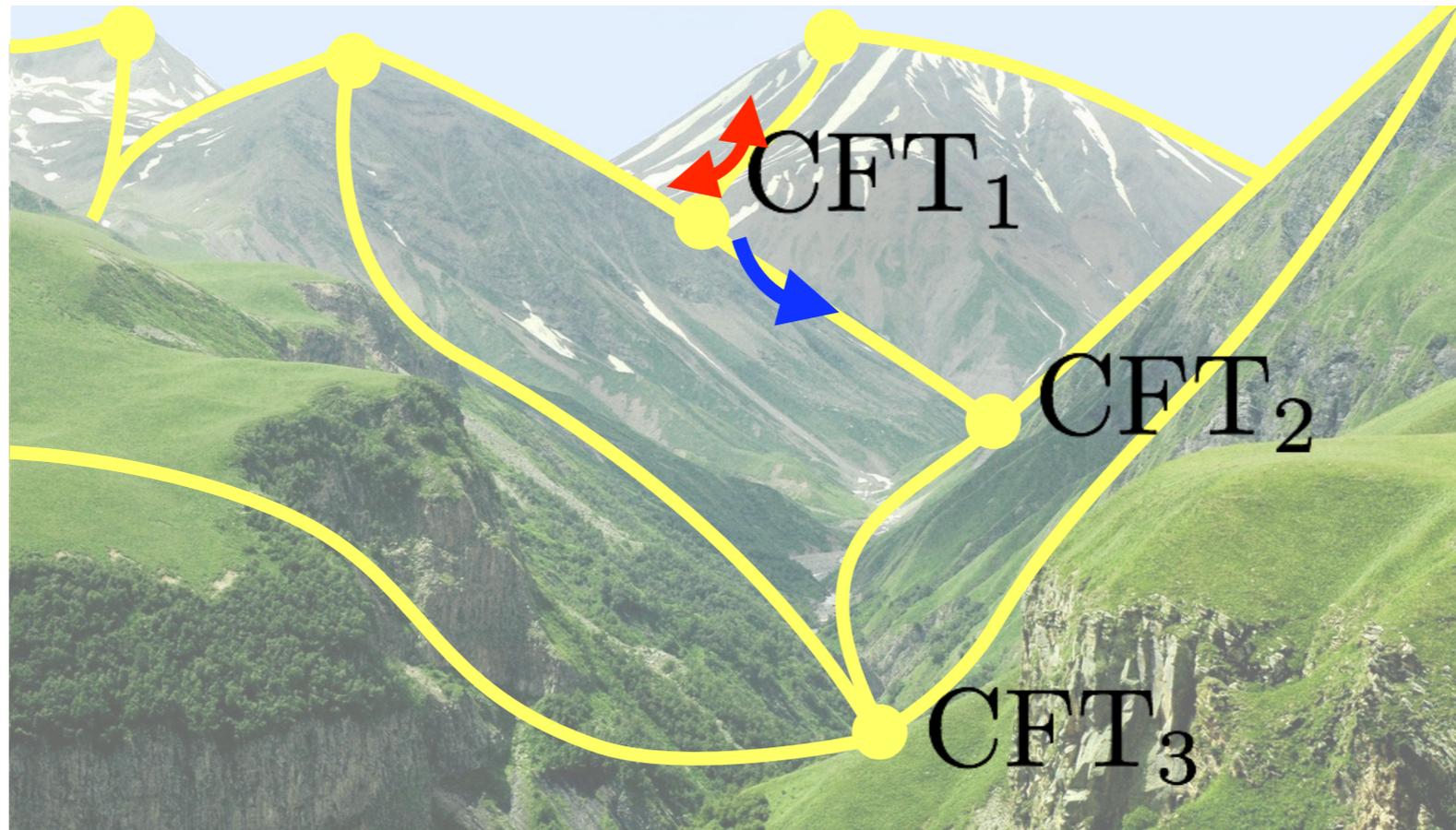


# RG Flow between CFTs



Add **irrelevant** term to CFT: RG flow pushes you back

# RG Flow between CFTs



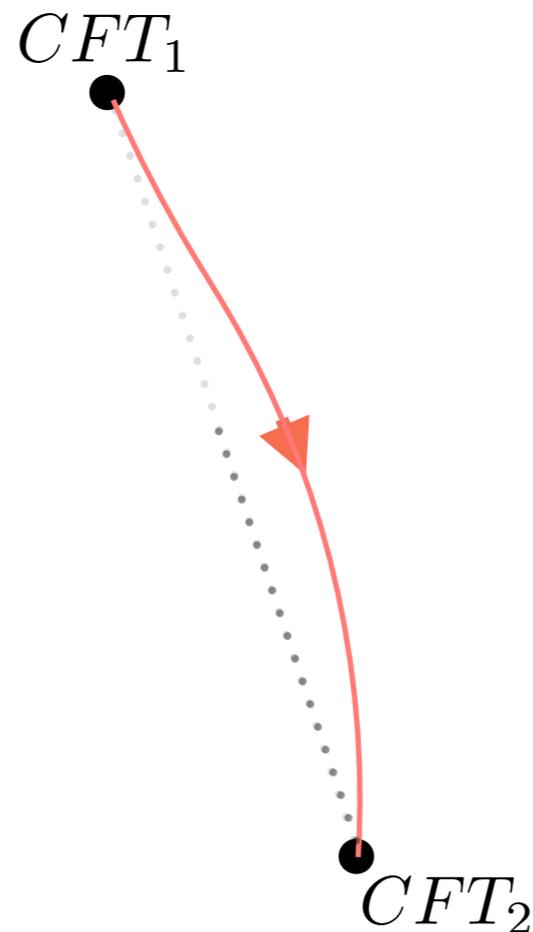
Add **irrelevant** term to CFT: RG flow  
pushes you back

Add *relevant* term to CFT: RG flow  
pushes you to new CFT

# Grand Bootstrap Goal: Solve QFT in 2 easy steps

Step 1: Solve CFT

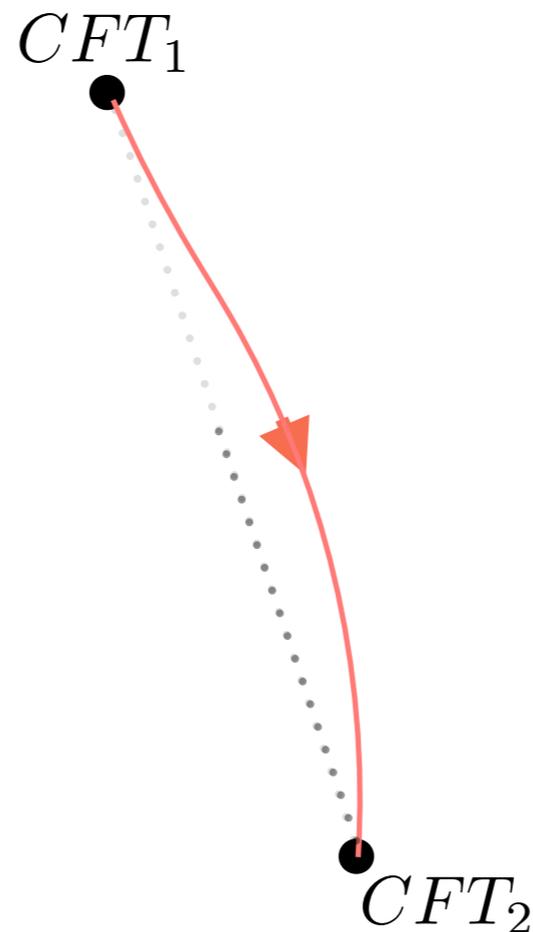
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# Grand Bootstrap Goal: Solve QFT in 2 easy steps

Step 1: Solve CFT

Step 2: Solve RG flow



# Conformal Bootstrap

1970's: Polyakov shows that form of 2- and 3-point functions in CFTs are fixed by symmetry

$$\langle \mathcal{O}_i(x_1) \mathcal{O}_j(x_2) \mathcal{O}_k(x_3) \rangle = \frac{c_{ijk}}{x_{12}^{\Delta_{ij,k}} x_{13}^{\Delta_{ik,j}} x_{23}^{\Delta_{jk,i}}}$$

and proposes bootstrap equation as consistency conditions on CFTs

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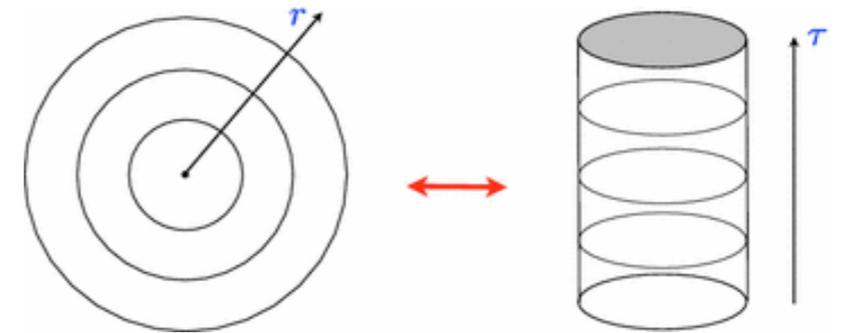
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$$\sum_{\alpha} \text{Tree}_1(\mathcal{O}_{\alpha}) = \sum_{\alpha} \text{Tree}_2(\mathcal{O}_{\alpha})$$

# Conformal Bootstrap

Key concepts behind bootstrap:

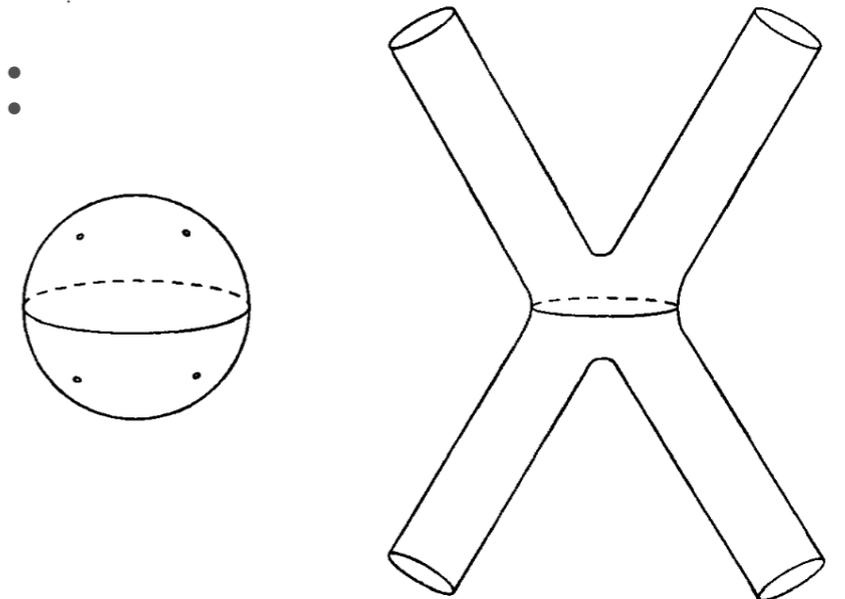
Radial quantization: conformal map from the plane to the (sphere) x (time)



Implies state-operator correspondence:

“states = local operators”

Infinite volume dynamics are controlled by finite volume system with a *discrete* set of “CFT” data



$$\{ \Delta_i, C_{ijkl} \}$$

# Conformal Bootstrap

Key concepts behind bootstrap:

**Operator Product Expansion:**  $\mathcal{O}_i \times \mathcal{O}_j = \sum_k c_{ijk} \mathcal{O}_k$

Start with insertion of two operators



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Decompose into a convenient basis at a fixed radius.

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Conformal Blocks

# Conformal Bootstrap

Key concepts behind bootstrap: **Bootstrap Equation**

Consider inserting 4 operators.

Powerful constraint: we can look at the wavefunction of the system on any surface and decompose into “conformal blocks” = “conformal spherical harmonics”

$$\begin{array}{cc} \rho(x_1) & \rho(x_4) \\ \times & \times \end{array}$$

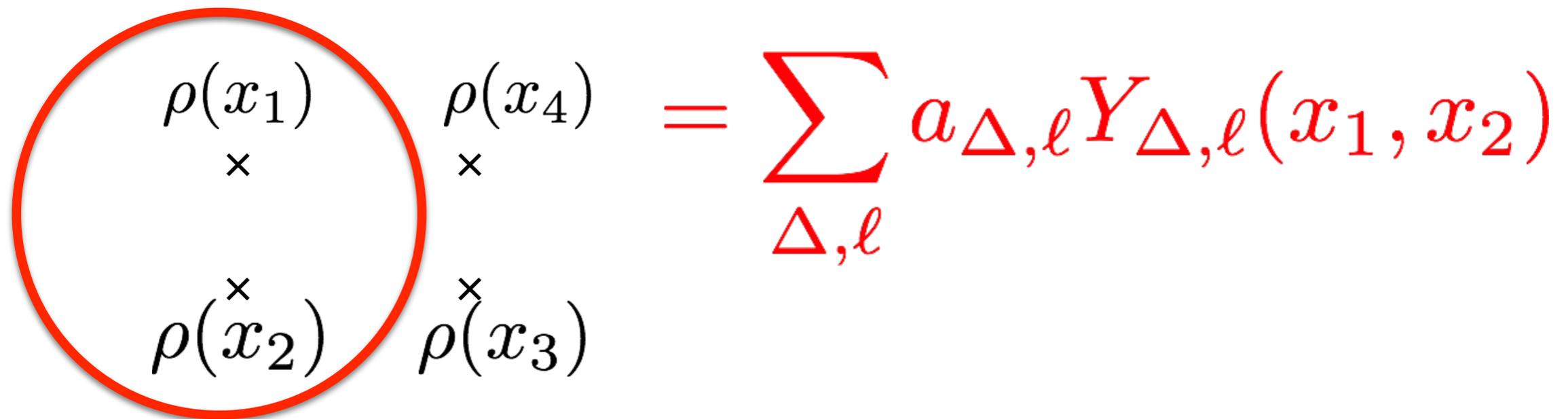
$$\begin{array}{cc} \rho^\times(x_2) & \rho^\times(x_3) \end{array}$$

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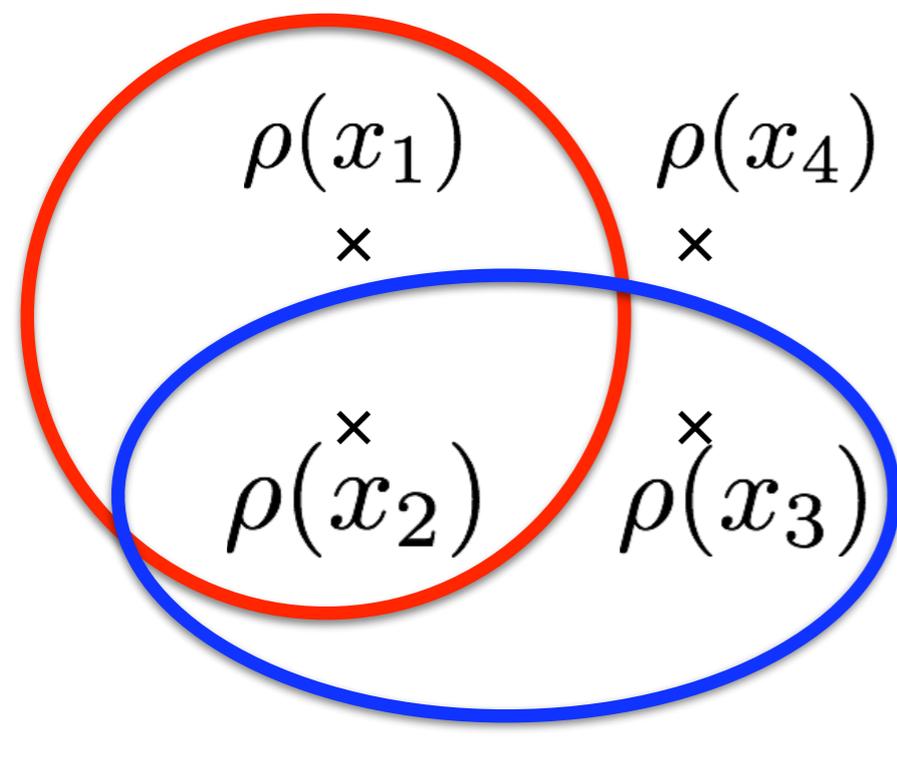

$$\begin{array}{cc} \rho(x_1) & \rho(x_4) \\ \times & \times \\ \rho(x_2) & \rho(x_3) \end{array} = \sum_{\Delta, \ell} a_{\Delta, \ell} Y_{\Delta, \ell}(x_1, x_2)$$

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$$\begin{aligned} \rho(x_1) \times \rho(x_4) &= \sum_{\Delta, \ell} a_{\Delta, \ell} Y_{\Delta, \ell}(x_1, x_2) \\ \rho(x_2) \times \rho(x_3) &= \sum_{\Delta, \ell} a_{\Delta, \ell} Y_{\Delta, \ell}(x_2, x_3) \end{aligned}$$

# Conformal Bootstrap

Key concepts behind bootstrap: **Bootstrap Equation**

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The diagram illustrates the bootstrap equation for a four-point correlation function. It features two conformal blocks separated by an equals sign, both summed over an index  $\alpha$ . The left block has a vertical line connecting two vertices, with an operator  $\mathcal{O}_\alpha$  on the right. The right block has a horizontal line connecting two vertices, with an operator  $\mathcal{O}_\alpha$  above it. Annotations include a red circle on the left, a blue circle on the left, a red bracket above the right block, a blue bracket below the right block, and labels  $\rho(\dots, x_2)$  and  $\rho(\dots, x_3)$  on the right. A blue label  $\Delta, l$  is at the bottom.

$$\sum_{\alpha} \text{[Block 1]} = \sum_{\alpha} \text{[Block 2]}$$

$\rho(\dots, x_2)$

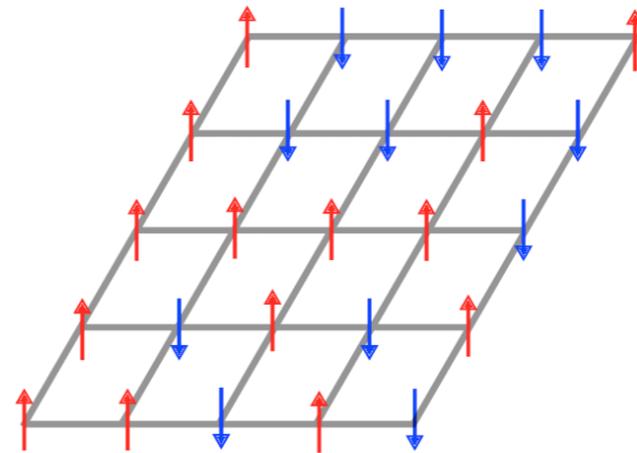
$\rho(\dots, x_3)$

$\Delta, l$

# Conformal Bootstrap

1980's: Belavin, Polyakov, & Zamolodchikov show how to use these ideas to solve special\* 2d CFTs

eg 2d Ising model



\*Crucially, these special CFTs have only a *finite* amount of “CFT data”  $\{\Delta_i, c_{ijk}\}$

But most 2d CFTs, and all CFTs at  $d > 2$ , have an infinite amount of CFT data

# Conformal Bootstrap Rebirth

2008: Rattazzi, Rychkov, Tonni, & Vichi

(Also, 2009: Simeon Kellerman - Modular Bootstrap)

Key idea: prove bounds by contradiction

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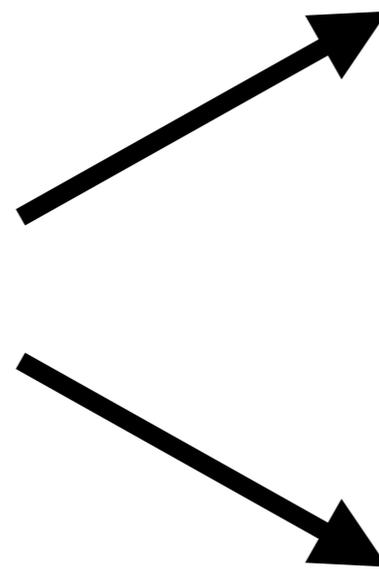
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*conformal  
oracle*

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*NO!*

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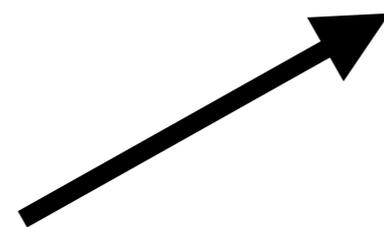
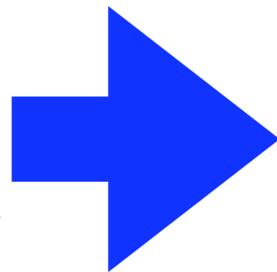
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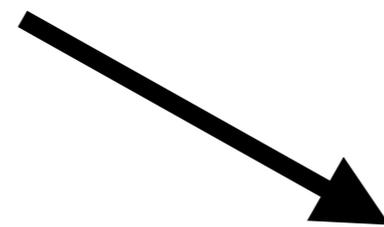
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*NO!*



*MAYBE!*

# Conformal Bootstrap Rebirth

## Conformal Oracle from Bootstrap Equation

First, do some trivial reorganization of bootstrap equation

$$\sum_{\Delta, \ell} a_{\Delta, \ell} Y_{\Delta, \ell}(x_1, x_2) = \sum_{\Delta, \ell} a_{\Delta, \ell} Y_{\Delta, \ell}(x_2, x_3)$$

# Conformal Bootstrap Rebirth

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$$\sum_{\Delta, \ell} a_{\Delta, \ell} Y_{\Delta, \ell}(x_1, x_2) = \sum_{\Delta, \ell} a_{\Delta, \ell} Y_{\Delta, \ell}(x_2, x_3)$$

Rewrite as

$$0 = \sum_{\Delta, \ell} a_{\Delta, \ell} \overbrace{Y_{\Delta, \ell}(x_1, x_2) - Y_{\Delta, \ell}(x_2, x_3)}^{F(x_i)}$$

# Conformal Bootstrap Rebirth

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Rewrite as

$$0 = \sum_{\Delta, \ell} a_{\Delta, \ell} \overbrace{F(x_i)}^{Y_{\Delta, \ell}(x_1, x_2) - Y_{\Delta, \ell}(x_2, x_3)} = F_{\text{vac}}(x_i) + \sum_{\Delta, \ell \neq \text{vac}} a_{\Delta, \ell} F_{\Delta, \ell}(x_i)$$

# Conformal Bootstrap Rebirth

Next, assume “partial CFT data”  $\{\Delta_i, c_{ijk}\}$  is correct, try to rule it out by proving a contradiction

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$$0 = F_{\text{vac}}(x_i) + \sum_{\Delta, \ell \neq \text{vac}} a_{\Delta, \ell} F_{\Delta, \ell}(x_i)$$

Idea: numerically find a linear functional  $\rho : F_{\Delta, \ell} \rightarrow \mathbb{R}$  such that

$$\rho(F_{\text{vac}}) = 1 \quad \text{and} \quad \rho(F_{\Delta, \ell}) \geq 0 \quad \text{for all } \Delta, \ell \text{ in } \{\Delta_i, c_{ijk}\}$$

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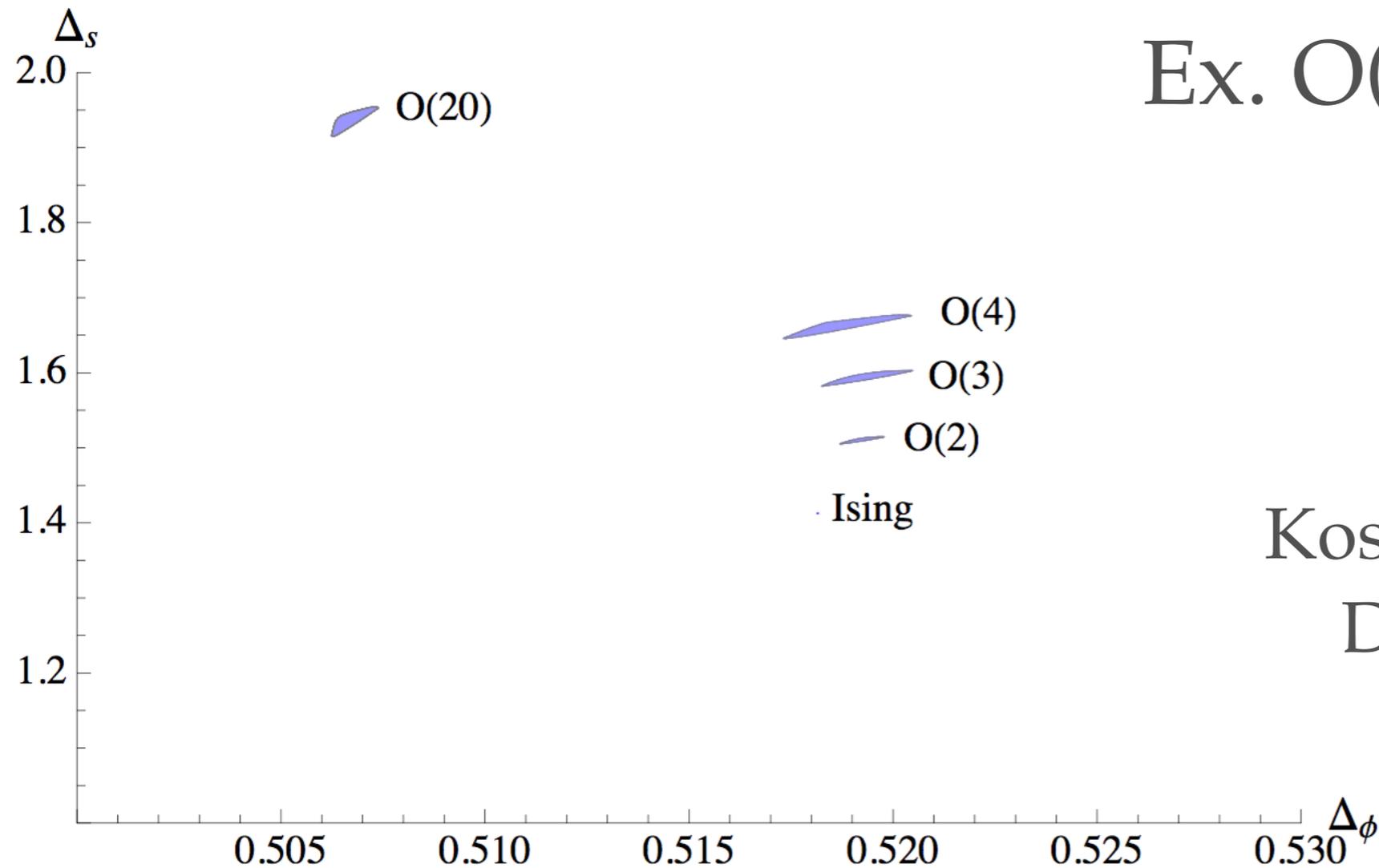
$$\rho(F_{\text{vac}}) = 1 \quad \text{and} \quad \rho(F_{\Delta, \ell}) \geq 0 \quad \text{for all } \Delta, \ell \text{ in } \{\Delta_i, c_{ijk}\}$$

If you can find such a functional, then, acting on the bootstrap equation, it implies

$$0 \geq 1 \quad \text{Contradiction!}$$

# 3d Islands

The  $O(N)$  archipelago



Ex.  $O(N)$  models in  
 $d=2+1$

Kos, Poland, Simmons-  
Duffin, Vichi, 2015

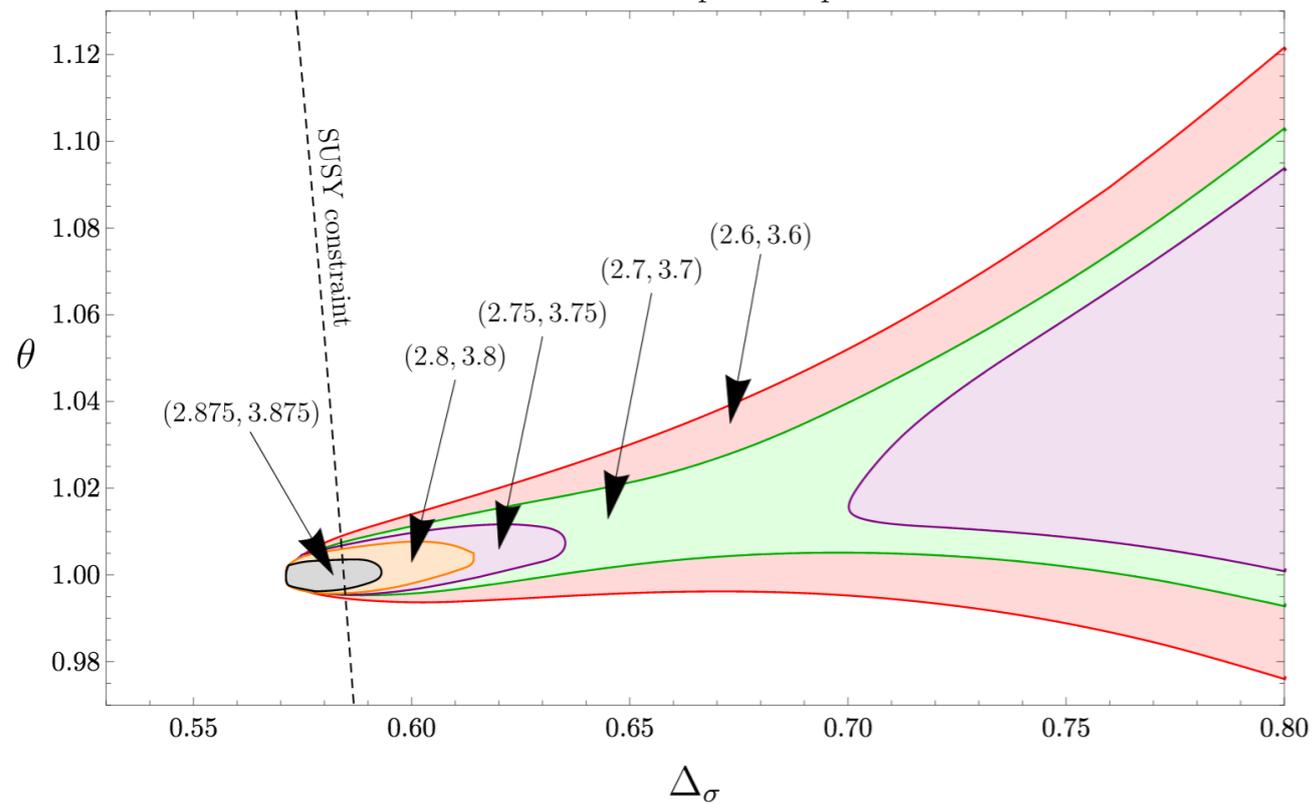
Ising “island” almost too small to see by eye

**At this point, can extract many OPE coefficients as well**

# 3d Islands

However, it's still a bit of an art to "carve out" islands

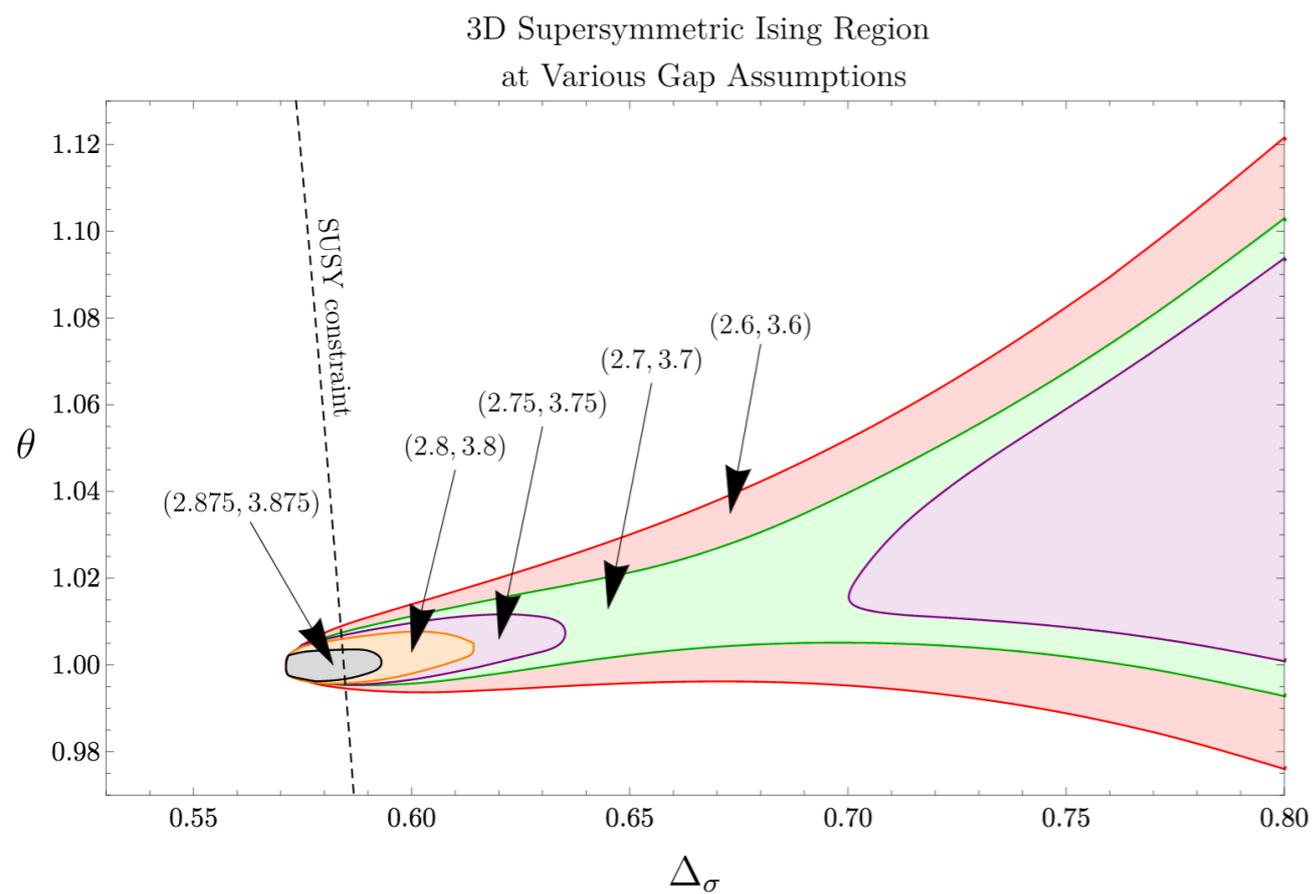
3D Supersymmetric Ising Region  
at Various Gap Assumptions



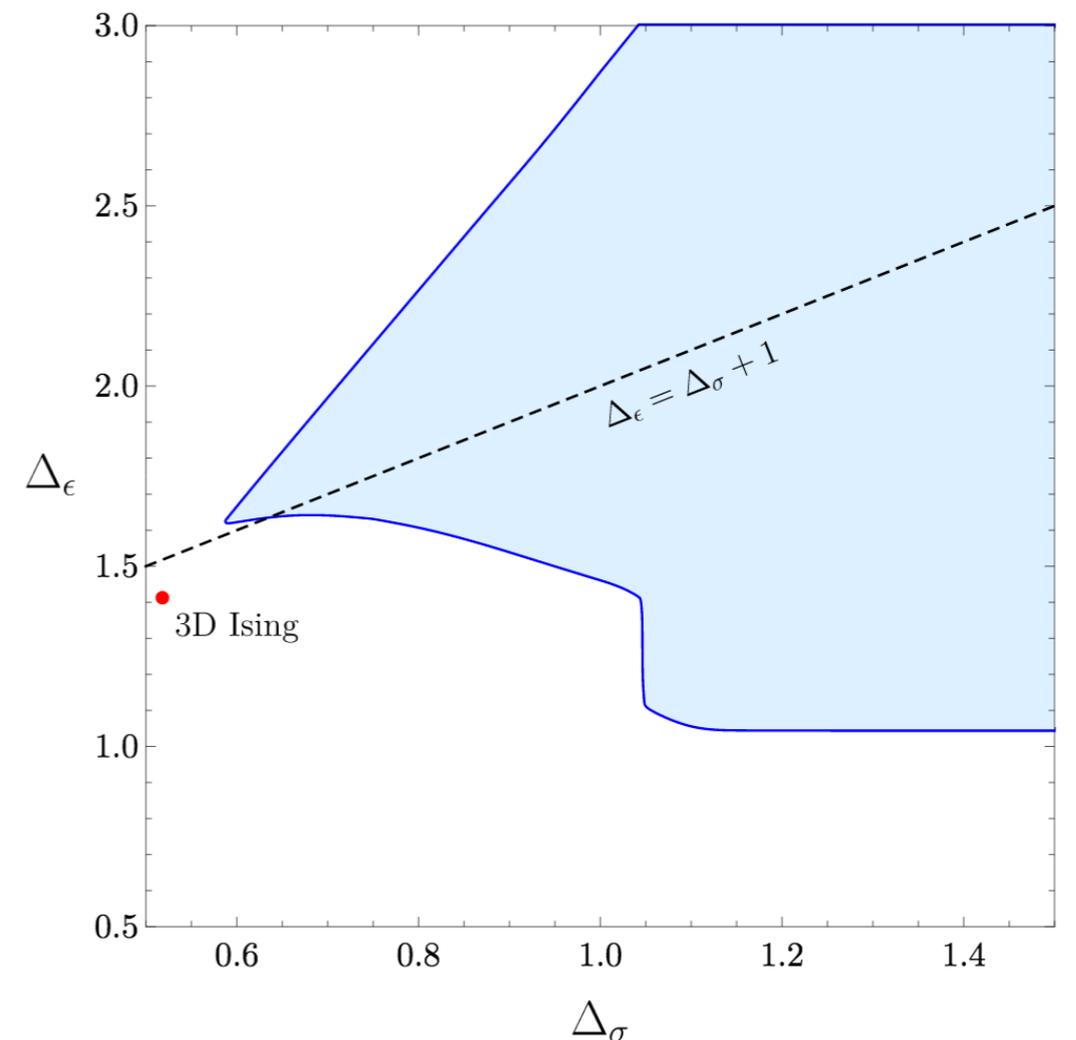
# 3d Islands

However, it's still a bit of an art to "carve out" islands

and not fully understood why it works well sometimes but not others



Allowed Region with Two Relevant Operators



# Grand Bootstrap Goal: Solve QFT in 2 easy steps

Step 1: Solve CFT

Step 2: Solve RG flow

$CFT_1$



$CFT_2$

# What about RG flows?

RG flow triggered by relevant deformation

$$H = H_{\text{CFT}} + \lambda \mathcal{O}_{\text{Relevant}}$$

*solved in "step 1"*

*"interaction" term*

# What about RG flows?

RG flow triggered by relevant deformation

$$H = H_{\text{CFT}} + \lambda \mathcal{O}_{\text{Relevant}}$$

*solved in "step 1"*      *"interaction" term*

Idea: Do *exact* diagonalization on  $H$  in a basis of eigenstates of  $H_{\text{CFT}}$  (Rayleigh-Ritz)

$$\langle \mathcal{O}_1 | H | \mathcal{O}_2 \rangle$$

Matrix elements just depend on UV CFT data

# Truncation Warm-up: Anharmonic Oscillator in QM

$$\mathcal{L} = \frac{1}{2}\dot{x}^2 - \frac{1}{2}m^2x^2 - \overbrace{gx^4}^V \longrightarrow H = H_{\text{SHO}} + V$$

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$V \sim g(a + a^\dagger)^4$

Basis:  $|n\rangle = \frac{1}{\sqrt{n!}}(a^\dagger)^n|0\rangle$        $\langle n|H|m\rangle = \delta_{nm}E_n^{(\text{SHO})} + \langle n|V|m\rangle$

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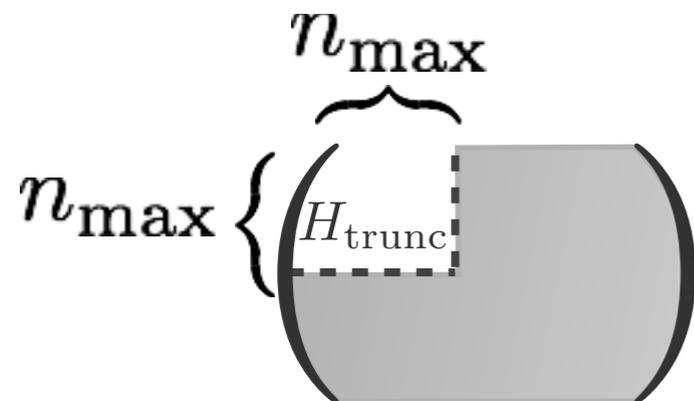
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Take finite  $n_{\text{max}}$  and diagonalize

$n_{\text{max}} \times n_{\text{max}}$  matrix

Eigenvalues converge quickly

as a function of  $n_{\text{max}}$



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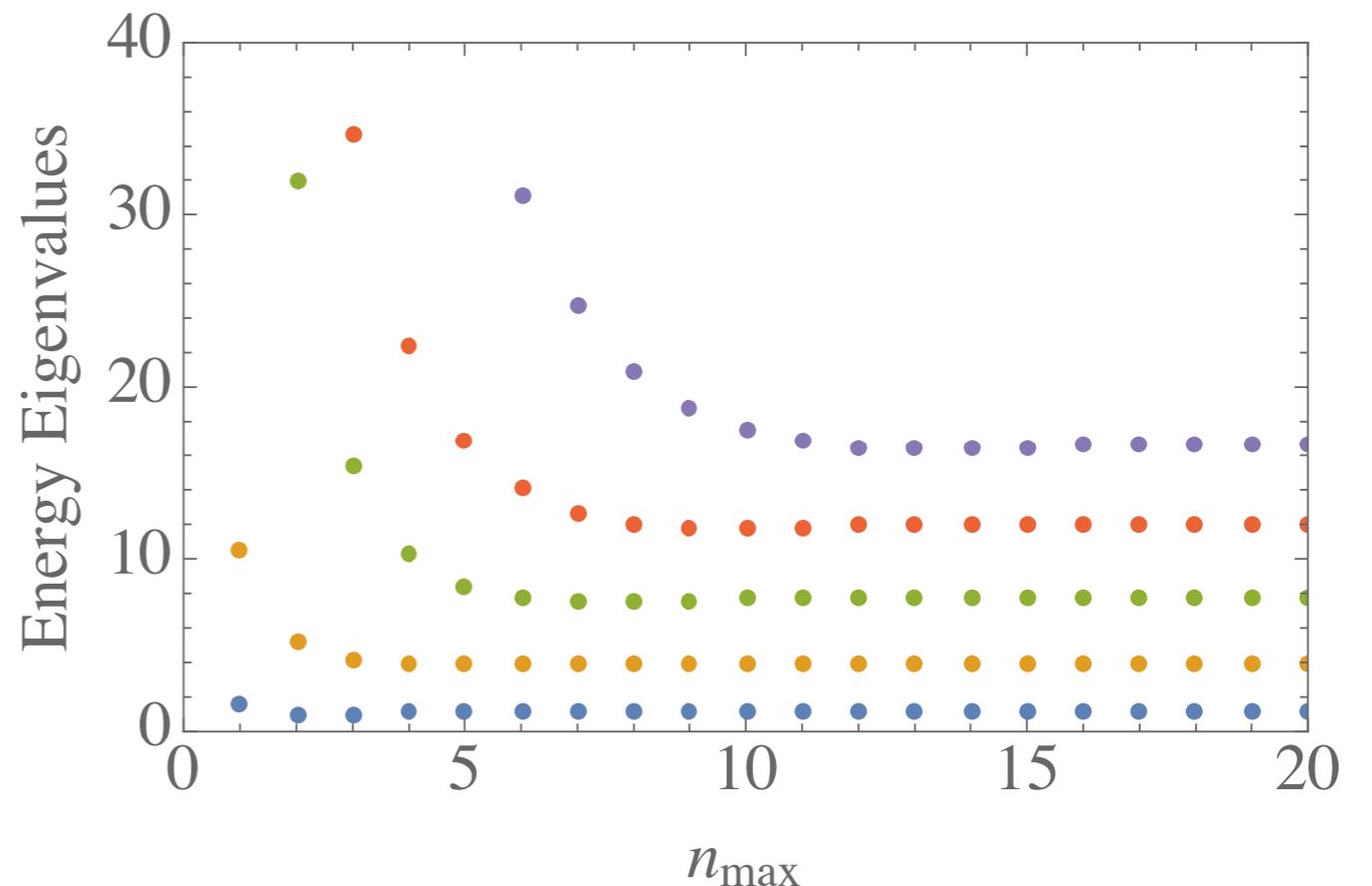
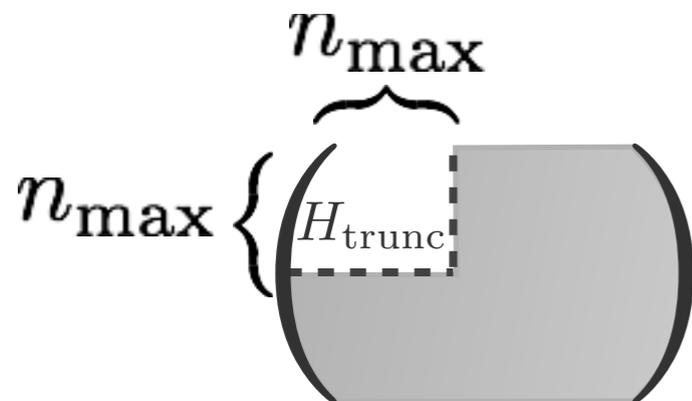
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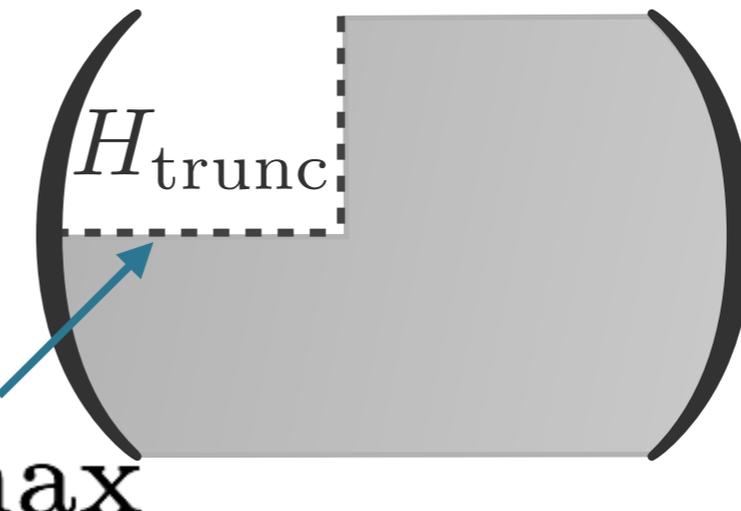
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as a function of  $n_{\text{max}}$



# Conformal Truncation

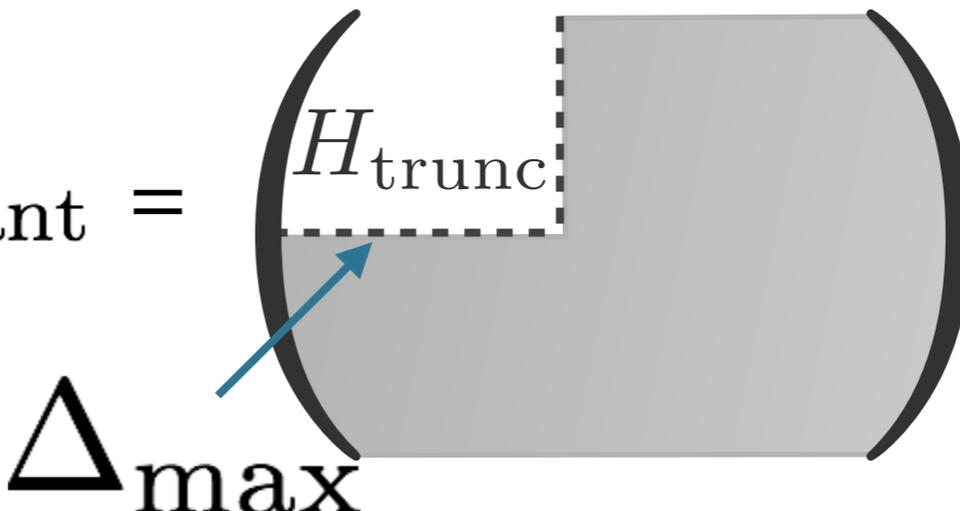
Same idea, applied to QFT  
Hamiltonian

$$H = H_{\text{CFT}} + \lambda \mathcal{O}_{\text{Relevant}} = \left( H_{\text{trunc}} \right)$$
A diagram illustrating conformal truncation. It shows a large gray semi-circular region bounded by a thick black arc on the right and a dashed line on the left. A horizontal dashed line extends from the center of the arc to the left, labeled  $H_{\text{trunc}}$ . A blue arrow points from the label  $\Delta_{\text{max}}$  below to the dashed line. The label  $\Delta_{\text{max}}$  is positioned below the dashed line, indicating the maximum dimension of states kept in the truncated theory.

Idea: keep only states with dimension below  $\Delta_{\text{max}}$

# Conformal Truncation

Same idea, applied to QFT  
Hamiltonian

$$H = H_{\text{CFT}} + \lambda \mathcal{O}_{\text{Relevant}} = \left( H_{\text{trunc}} \right)$$
A diagram illustrating conformal truncation. It shows a large gray semi-circular region representing the full Hamiltonian. A horizontal dashed line is drawn across the middle of the region, labeled  $H_{\text{trunc}}$ . A blue arrow points from the label  $\Delta_{\text{max}}$  below to the dashed line, indicating that the truncation is performed at a maximum dimension  $\Delta_{\text{max}}$ .

Idea: keep only states with dimension below  $\Delta_{\text{max}}$

$\langle \mathcal{O}_1 | H | \mathcal{O}_2 \rangle$  Matrix elements just depend on UV CFT data

$\langle \mathcal{O}_1 | H_{\text{CFT}} | \mathcal{O}_2 \rangle$  determined by CFT dimensions  $\Delta_i$

$\langle \mathcal{O}_1 | \mathcal{O}_{\text{Relevant}} | \mathcal{O}_2 \rangle$  determined by CFT OPE coefficients  $c_{ijk}$

# Conformal Truncation

1990 : Zamolodchikov & Yurov apply conformal truncation to deformations of solvable 2d CFTs in finite volume

More recent work:

$\phi^4$  theory in 2d  
2d + 3d

finite volume:

Hogervorst, Rychkov, Van Rees, Vitale '14

infinite volume:

Anand, Genest, Katz, Khandker, Walters '17

2d QCD @ finite  $N_c$

Katz, Tavares, Xu '14

3d U(1) Chern-Simons  
gauge theory at large  $N_f$

ALF, Katz, Vitale, Walters '18

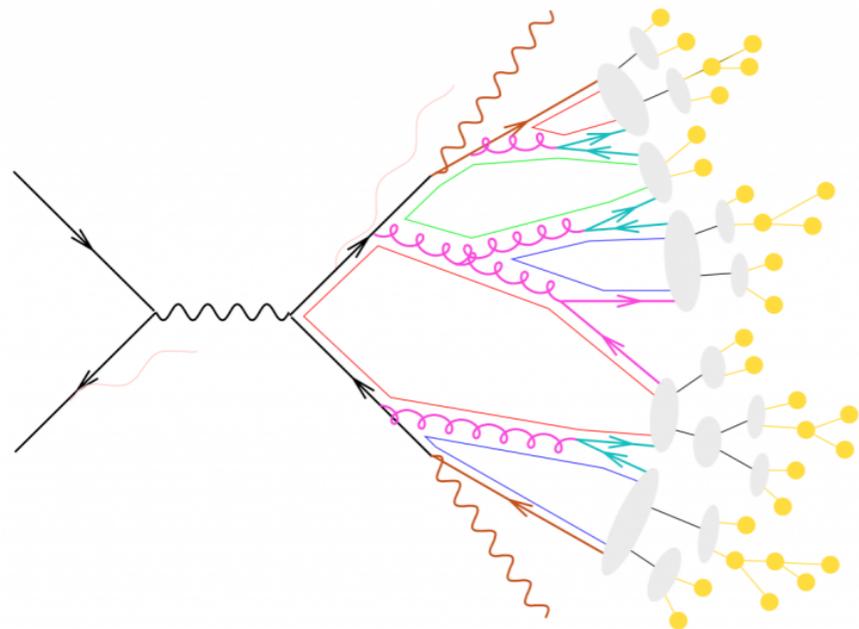
# Conformal Truncation

Because you diagonalize the Hamiltonian, you get eigenvalues *and* eigenvectors

This allows one to compute spectral functions, which encode correlation functions

$$\langle J^\mu(p) J^\nu(-p) \rangle = (q^2 \eta^{\mu\nu} - q^\mu q^\nu) \int_0^\infty ds \frac{\rho(s)}{p^2 - s + i\epsilon}$$

and scattering amplitudes

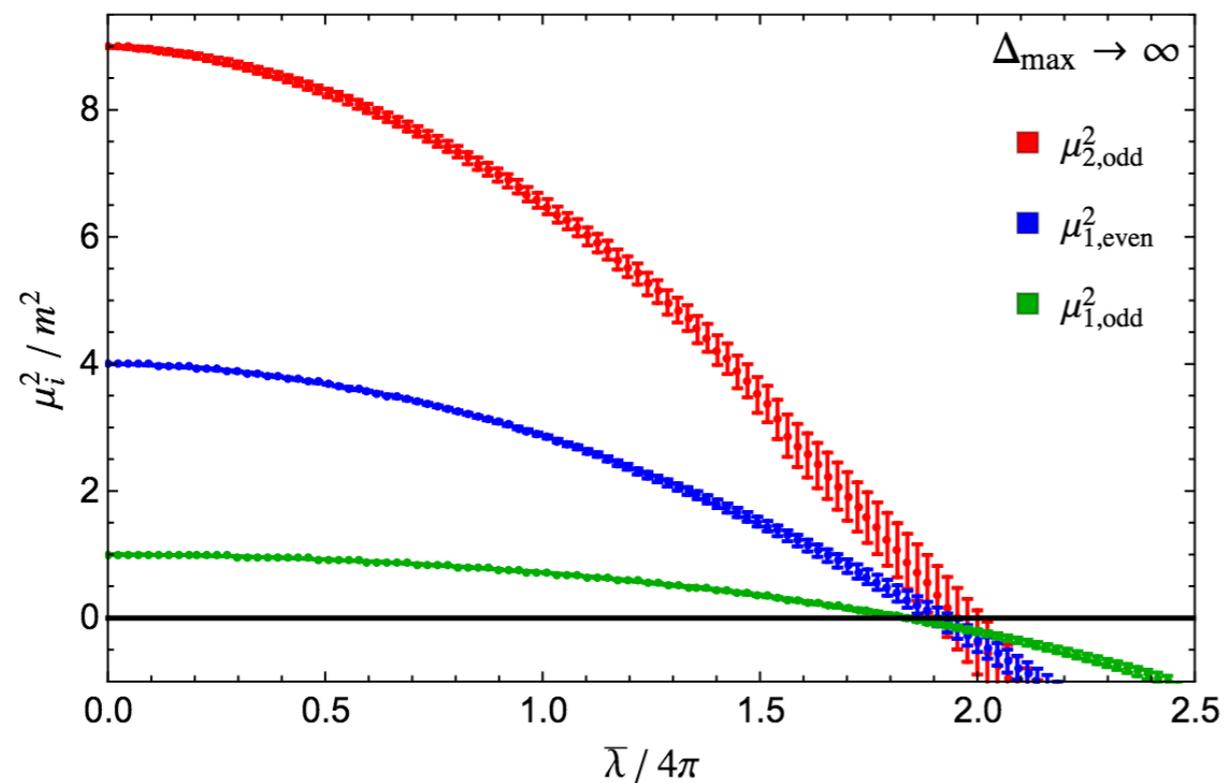


$$\sigma(e^+e^- \rightarrow \text{hadrons}) = -\frac{(4\pi\alpha)^2}{s} \rho(s)$$

# Conformal Truncation

Example  $d = 2$   $\mathcal{L} = \frac{1}{2}(\partial\phi)^2 - \frac{1}{2}m^2\phi^2 - \frac{\lambda}{4!}\phi^4$   $\bar{\lambda} = \frac{\lambda}{m^2}$

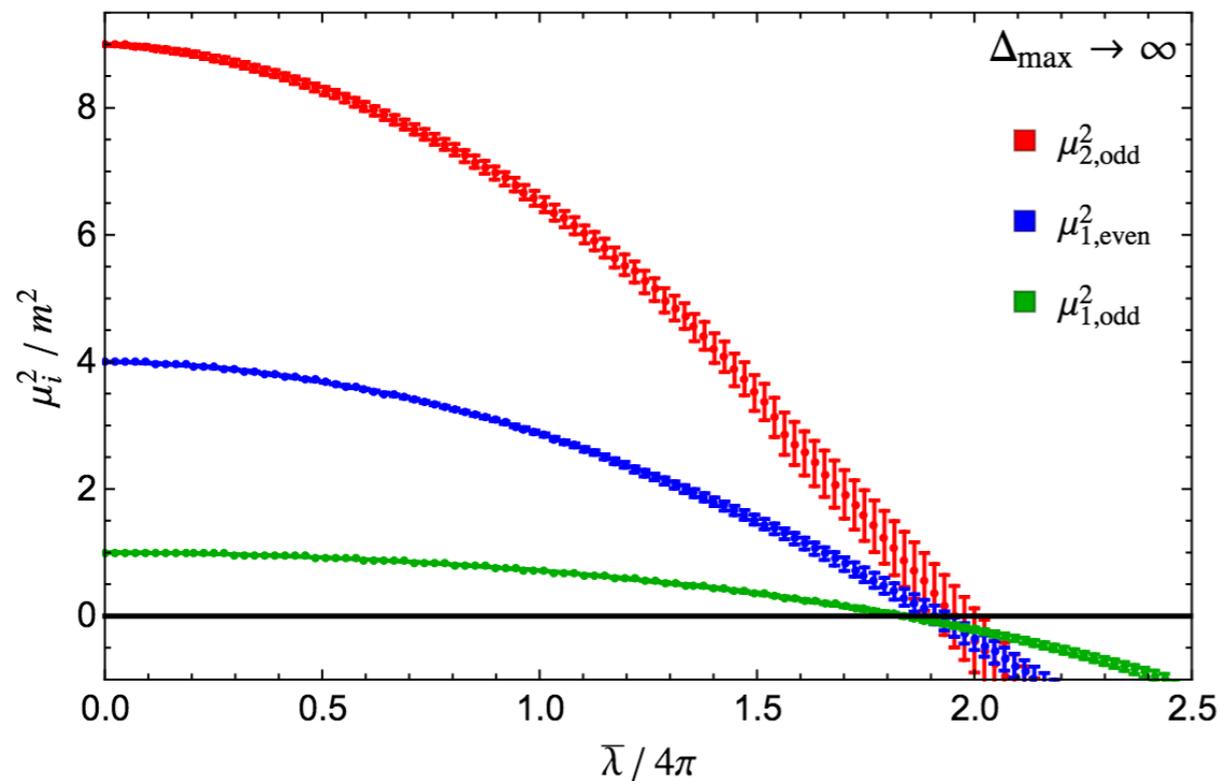
Eigenvalues:  
mass spectrum as a  
function of coupling



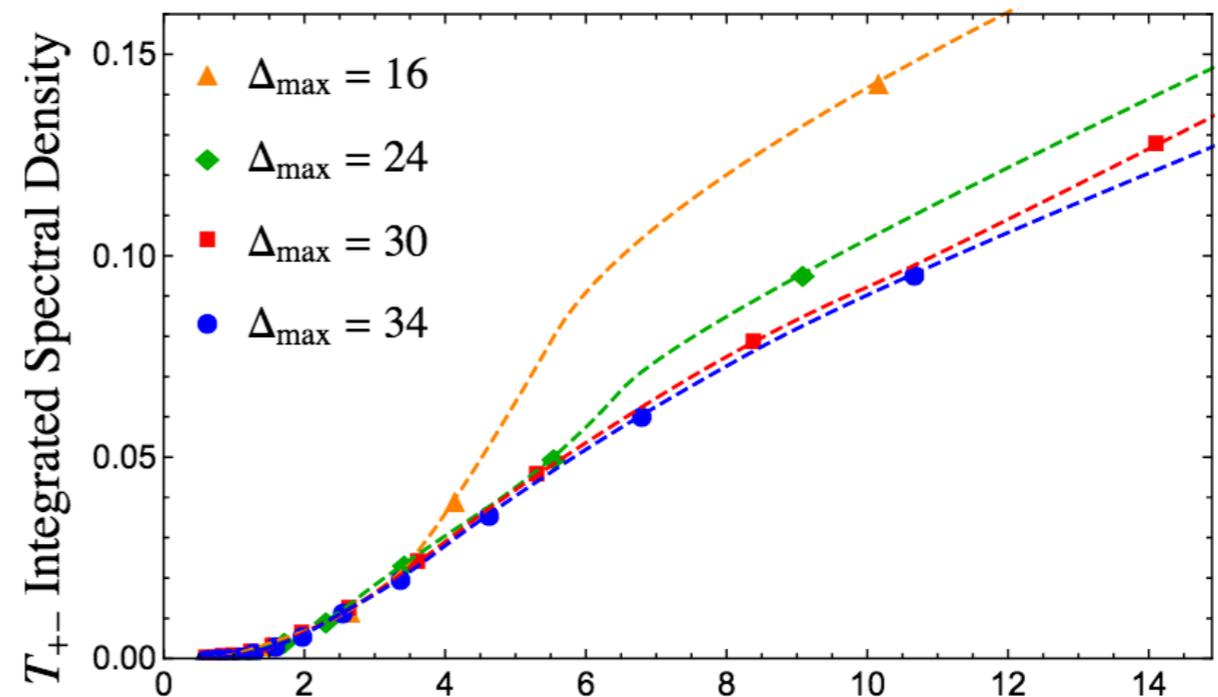
# Conformal Truncation

Example  $d = 2$   $\mathcal{L} = \frac{1}{2}(\partial\phi)^2 - \frac{1}{2}m^2\phi^2 - \frac{\lambda}{4!}\phi^4$   $\bar{\lambda} = \frac{\lambda}{m^2}$

Eigenvalues:  
mass spectrum as a  
function of coupling



Eigenvectors:  
spectral function of stress  
tensors as a function of  
momentum scale



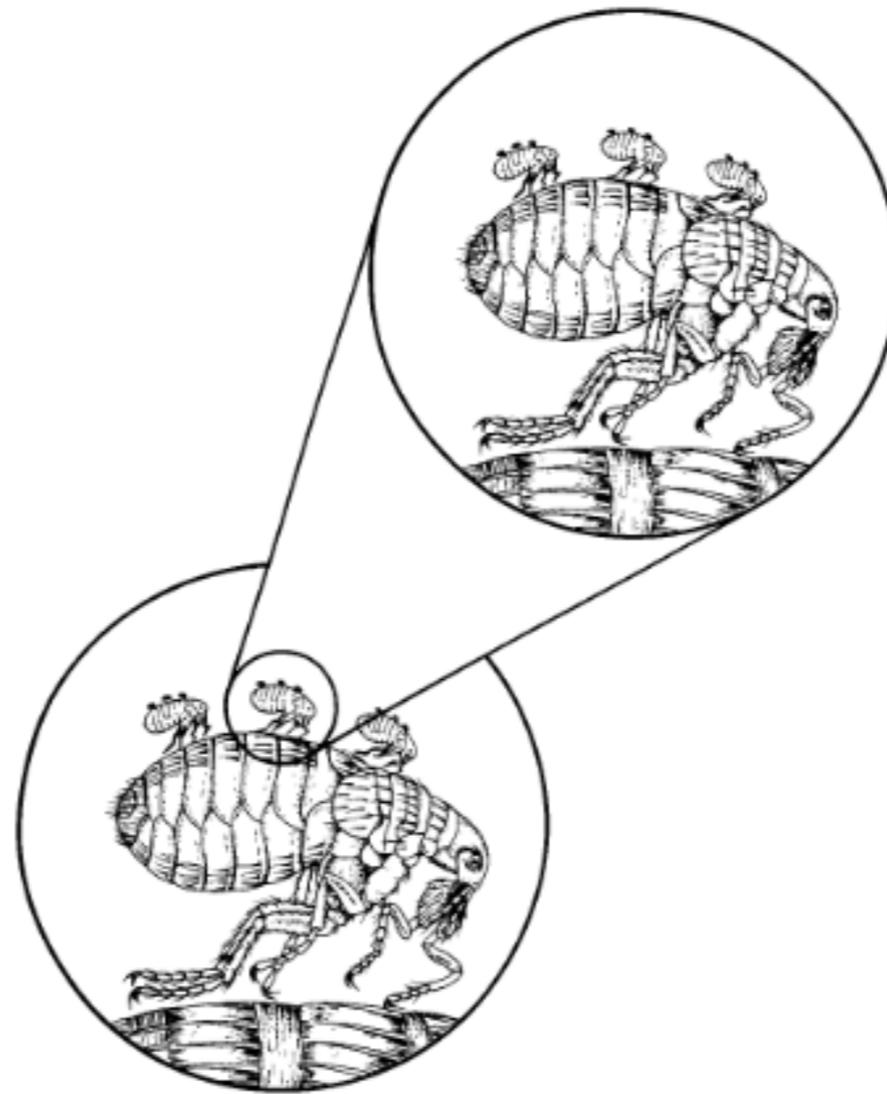
Anand et al. '17

$$\mathcal{I}_{T_{+-}}(\mu) \equiv \int_0^\mu d\mu'^2 \rho_{T_{+-}}(\mu')$$

# But many challenges remain...

- 1) At present, bootstrap cannot determine CFT data for most CFTs - only some “lucky” CFTs that happen to be at the edge of numeric bounds
- 2) Conformal truncation requires one to solve the UV CFT - can only use special solvable CFTs (e.g. free theories & integrable models)
- 3) Requires intensive numerical work - can more efficient methods be invented?
- 4) Etc.

# THE END



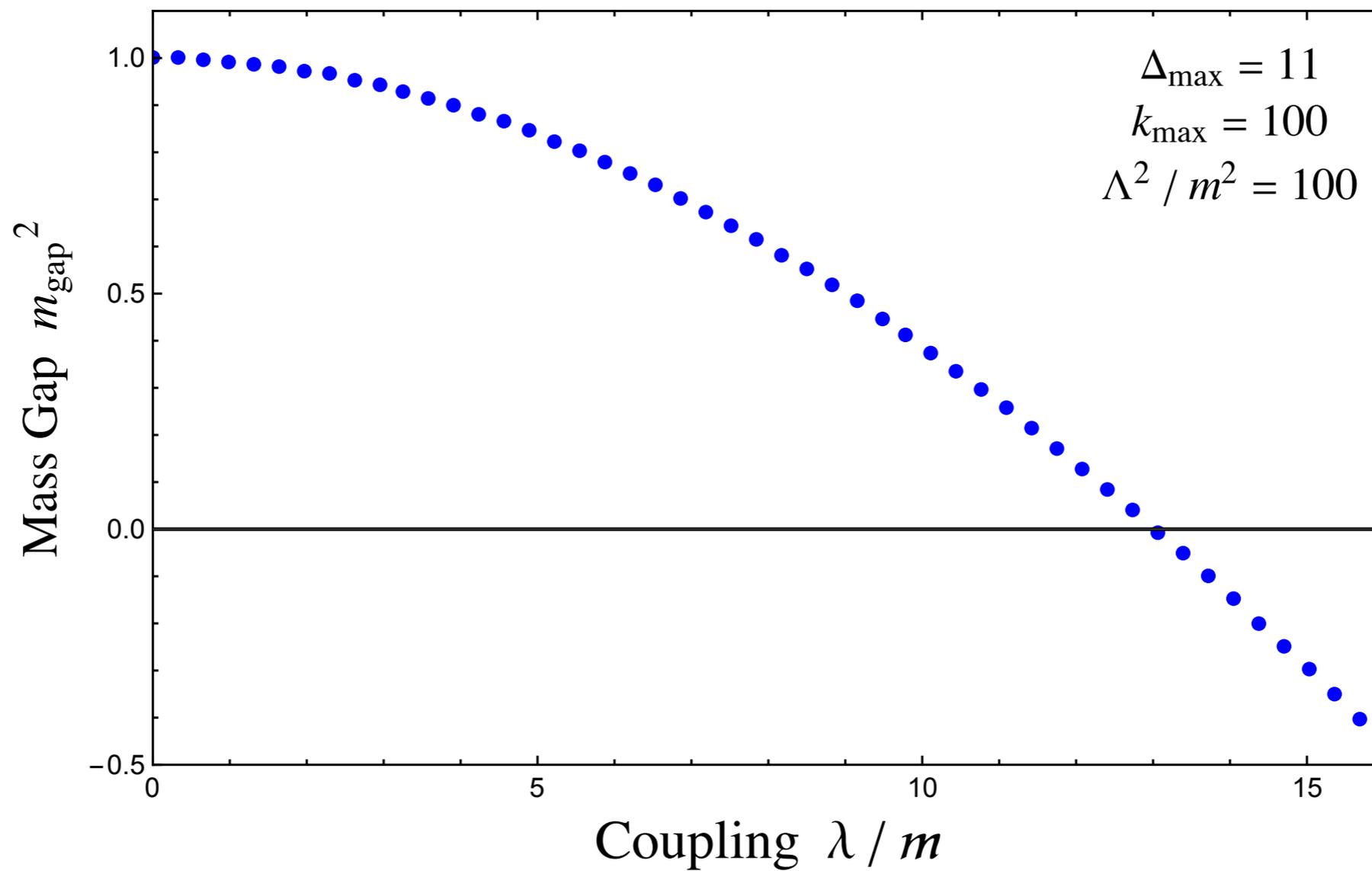


# Backup Slides

$$d = 3 \quad \mathcal{L} = \frac{1}{2}(\partial\phi)^2 - \frac{1}{2}m^2\phi^2 - \frac{\lambda}{4!}\phi^4$$

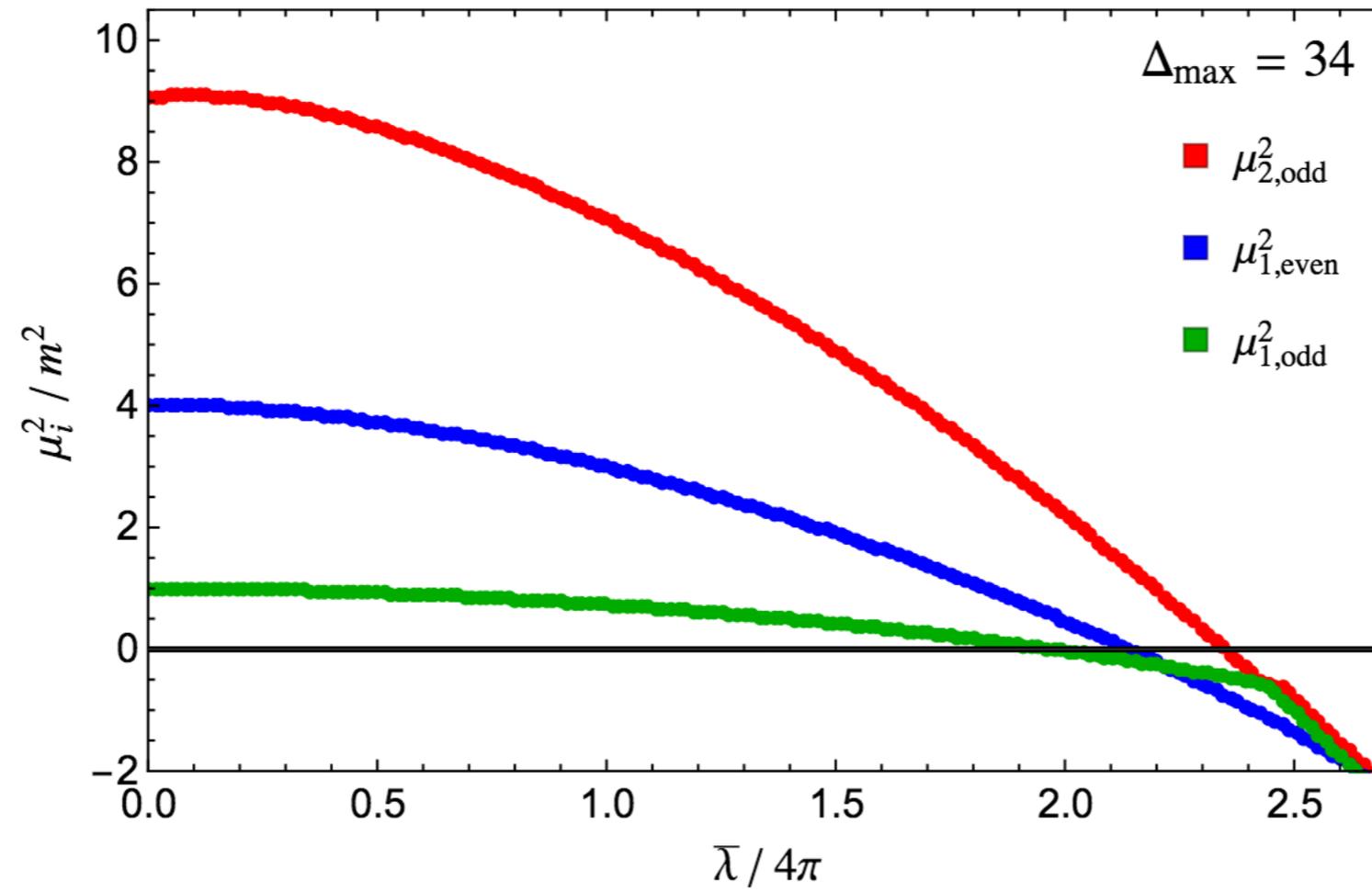
3d Ising

Preliminary



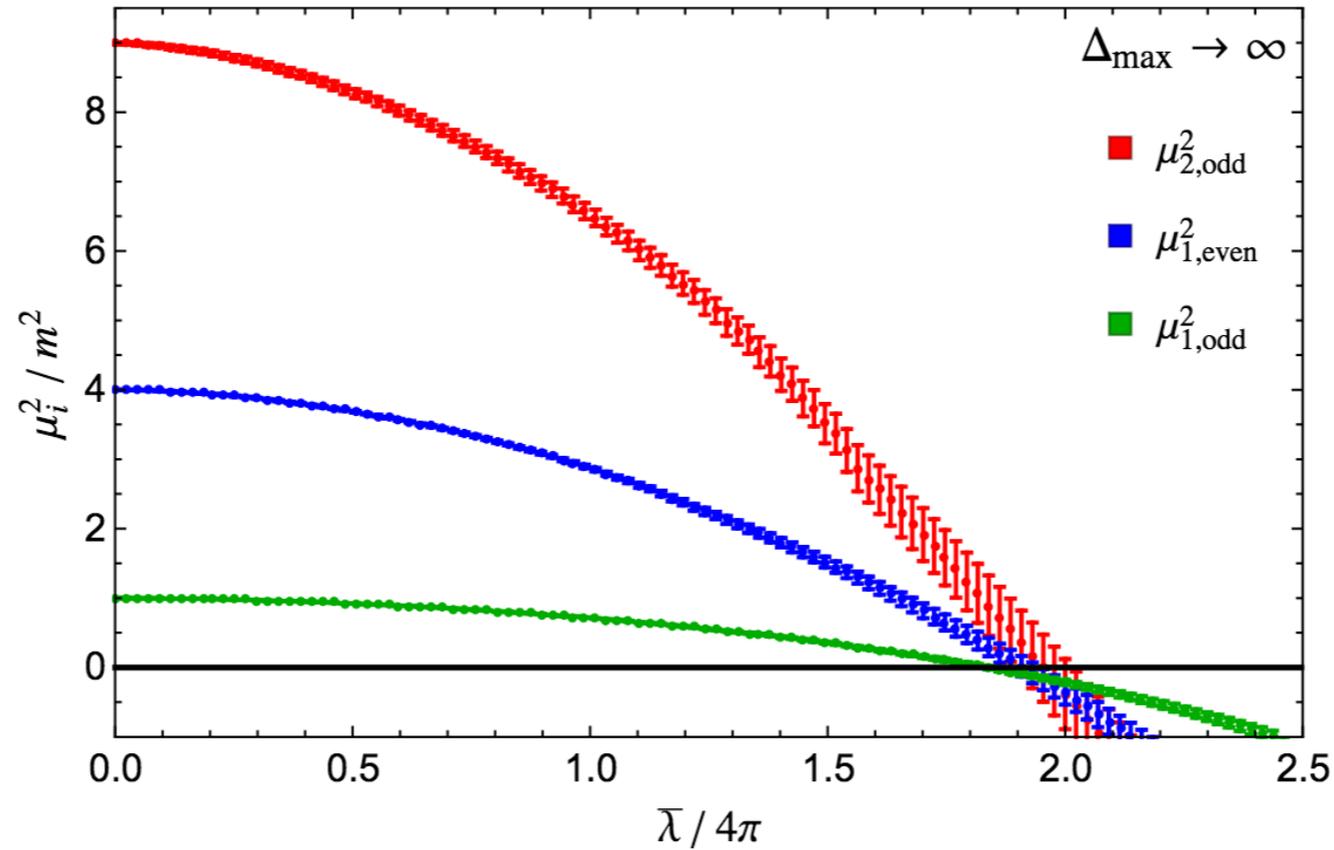
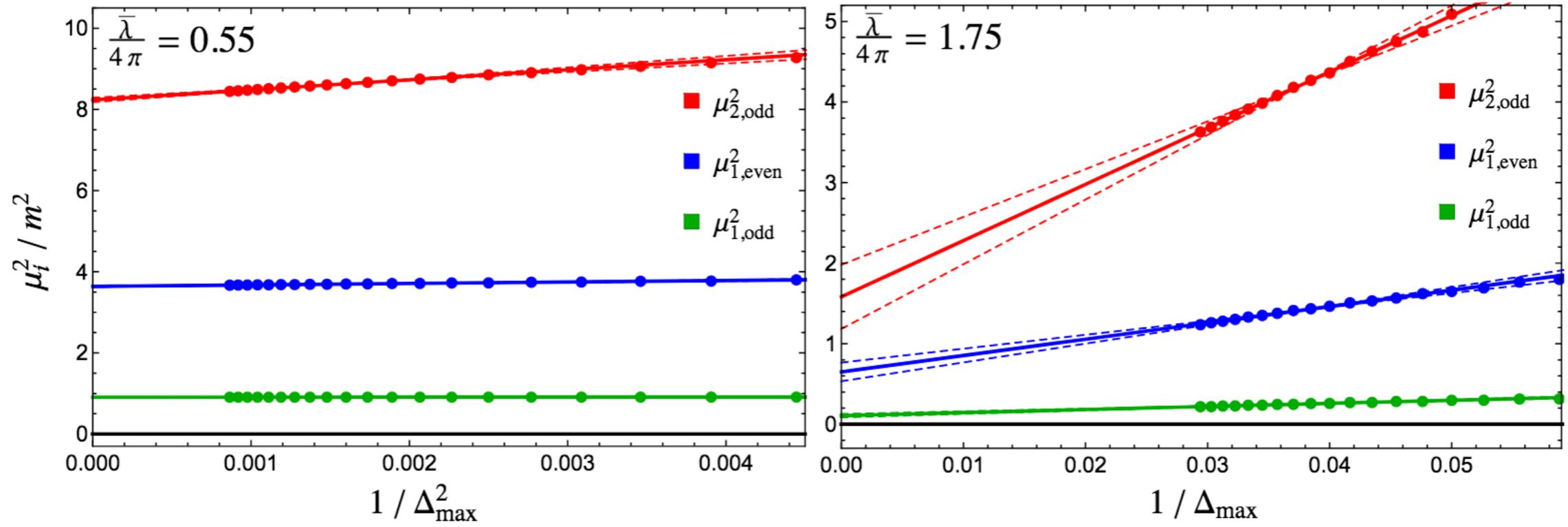
M. Walters

$$d = 2 \quad \mathcal{L} = \frac{1}{2}(\partial\phi)^2 - \frac{1}{2}m^2\phi^2 - \frac{\lambda}{4!}\phi^4 \quad \bar{\lambda} = \frac{\lambda}{m^2}$$



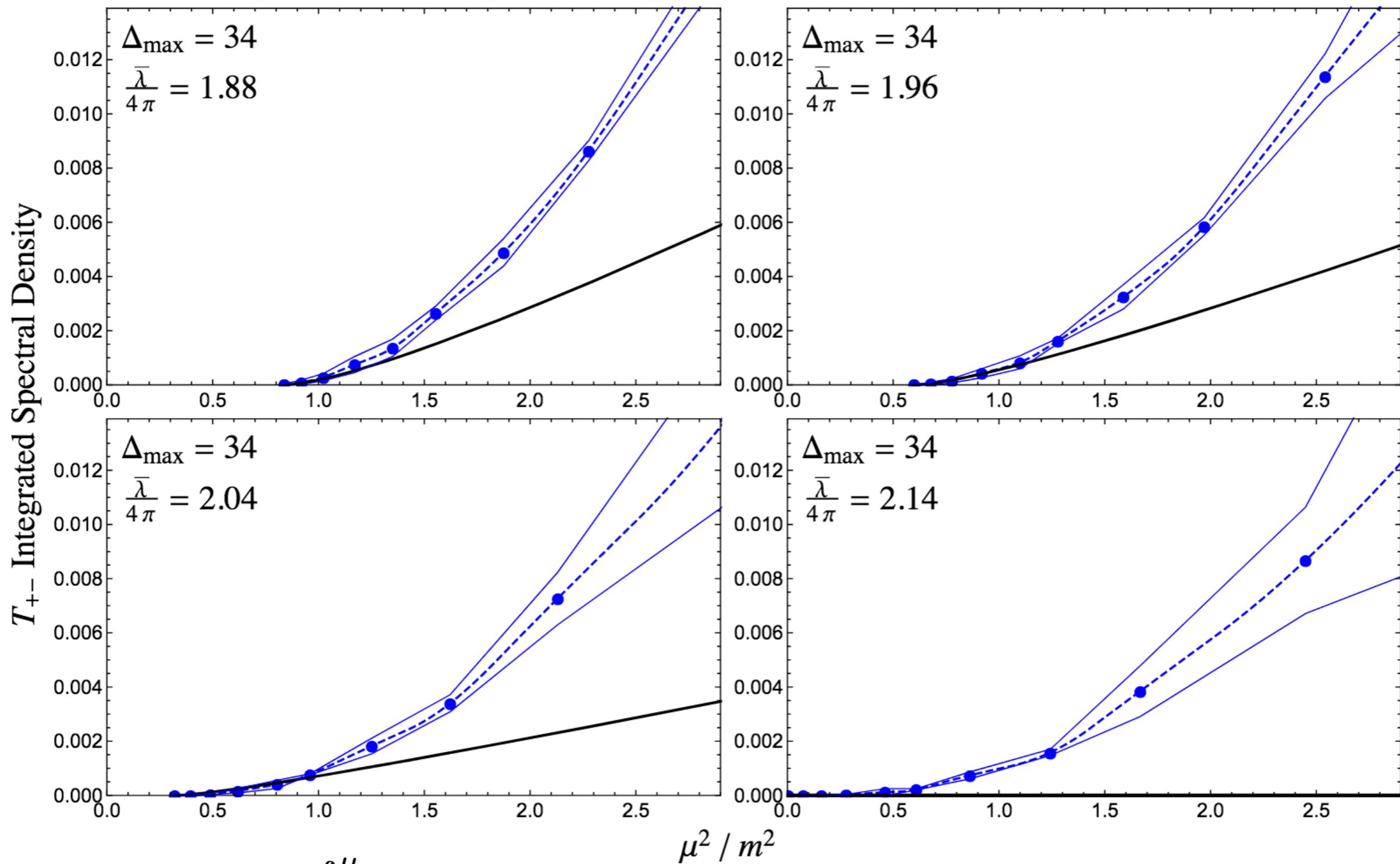
Anand et al. '17

$$d = 2 \quad \mathcal{L} = \frac{1}{2}(\partial\phi)^2 - \frac{1}{2}m^2\phi^2 - \frac{\lambda}{4!}\phi^4 \quad \bar{\lambda} = \frac{\lambda}{m^2}$$



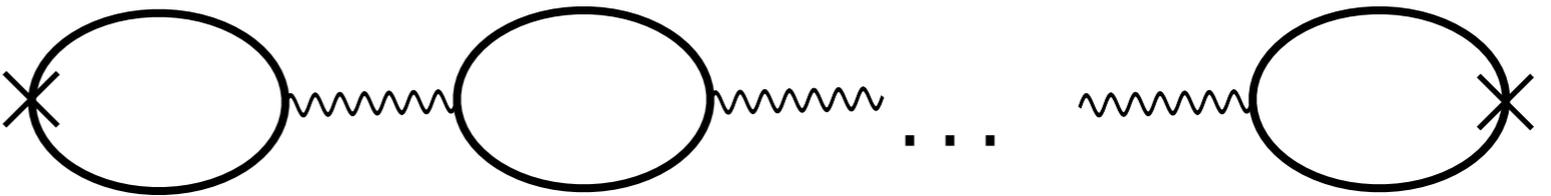
Anand et al. '17

$$d = 2 \quad \mathcal{L} = \frac{1}{2}(\partial\phi)^2 - \frac{1}{2}m^2\phi^2 - \frac{\lambda}{4!}\phi^4$$

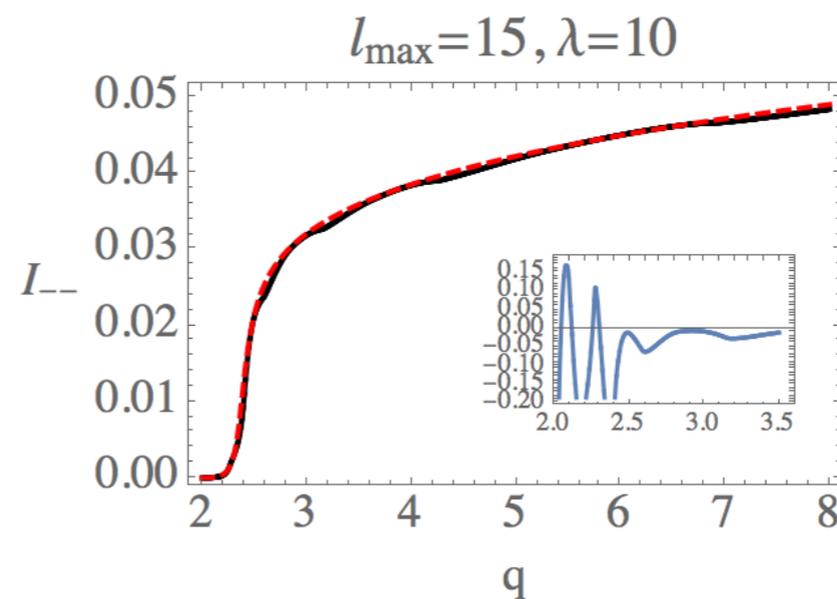
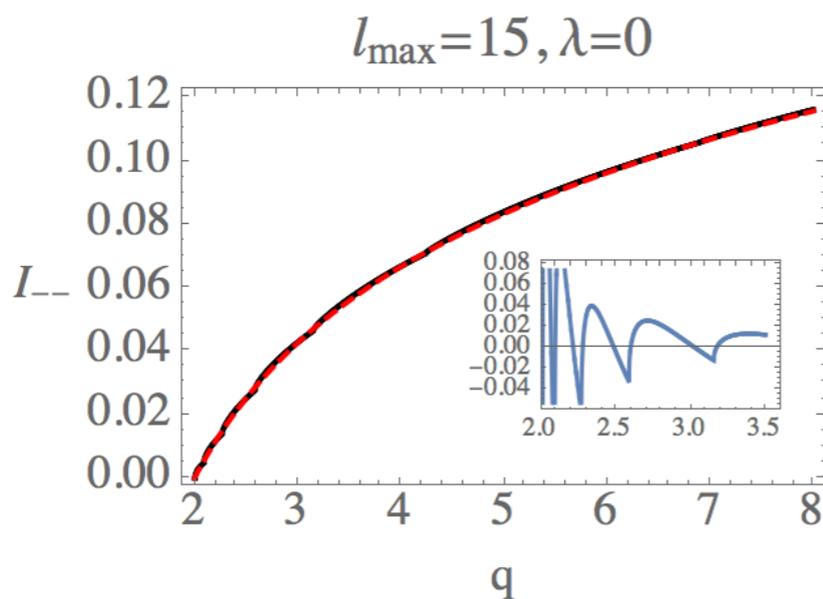
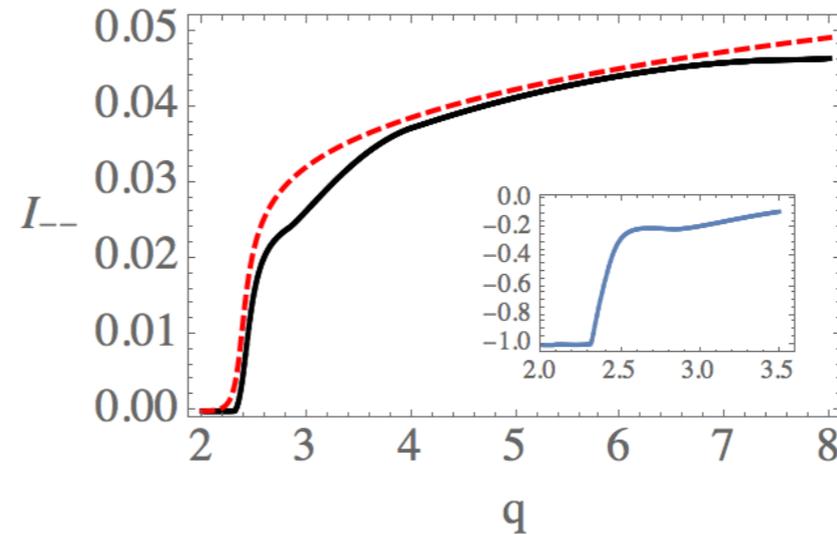
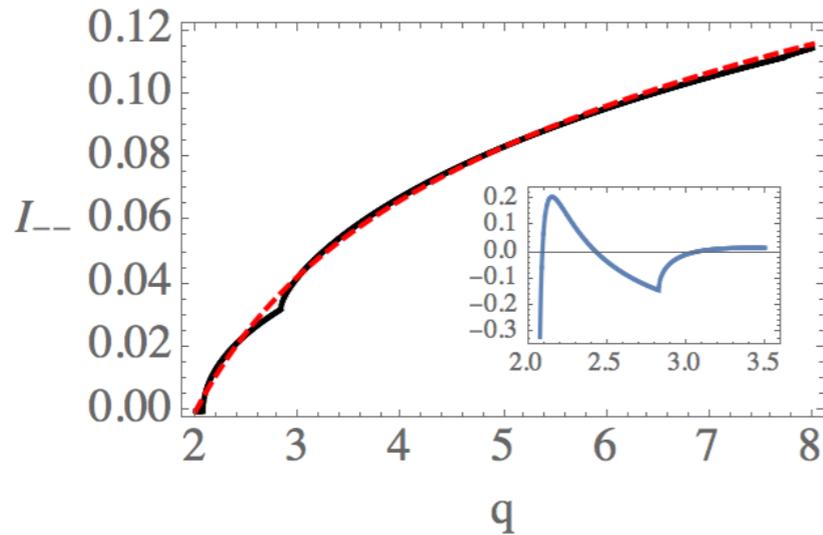


$$\mathcal{I}_{T_{+-}}(\mu) \equiv \int_0^\mu d\mu'^2 \rho_{T_{+-}}(\mu')$$

3d Chern-Simons  $N_f = \infty$   $\mathcal{L} = \bar{\Psi}(i\gamma^\mu D_\mu - m)\Psi + \frac{k}{4\pi}ada - m_a a^2$

$\langle J_\mu J_\nu \rangle =$  

$\pi\rho_{--}(q) = \text{Im} \left( \frac{32i\pi q ((q^2 + 4) \tanh^{-1}(\frac{q}{2}) - 2q)}{4q^2(8\pi q - \lambda)(\lambda + 8\pi q) + \lambda \tanh^{-1}(\frac{q}{2}) (-128\pi q^3 + 4\lambda(q^2 + 4)q - \lambda(q^2 - 4)^2 \tanh^{-1}(\frac{q}{2}))} \right)$



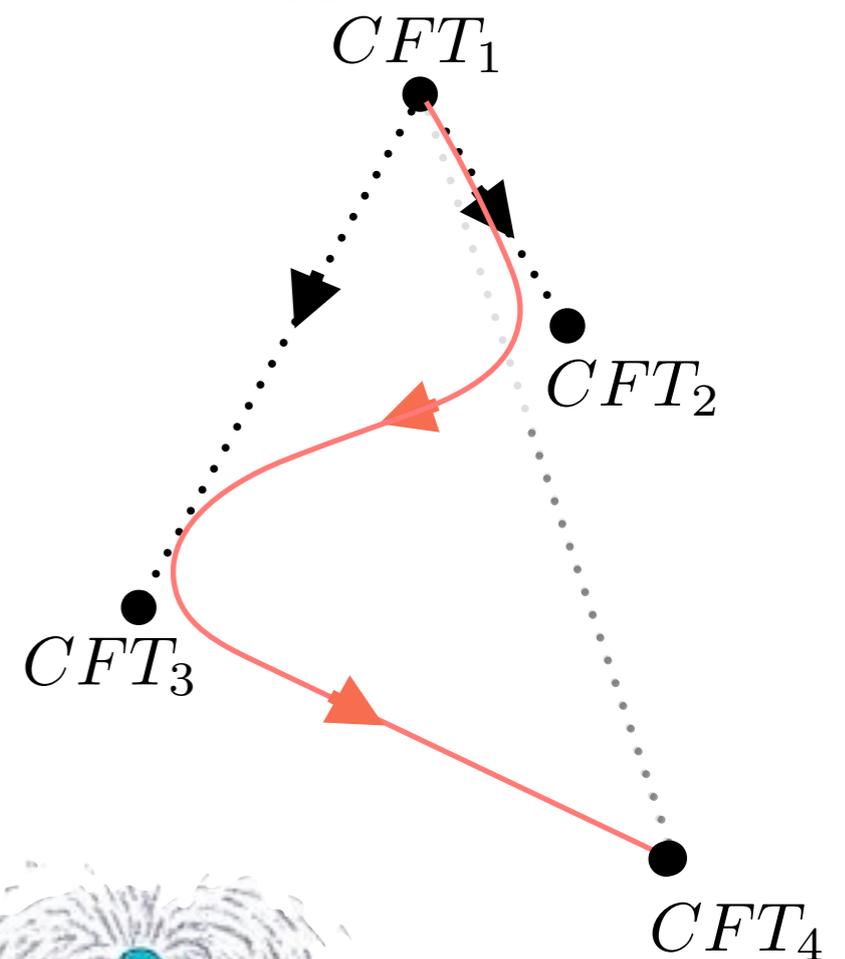
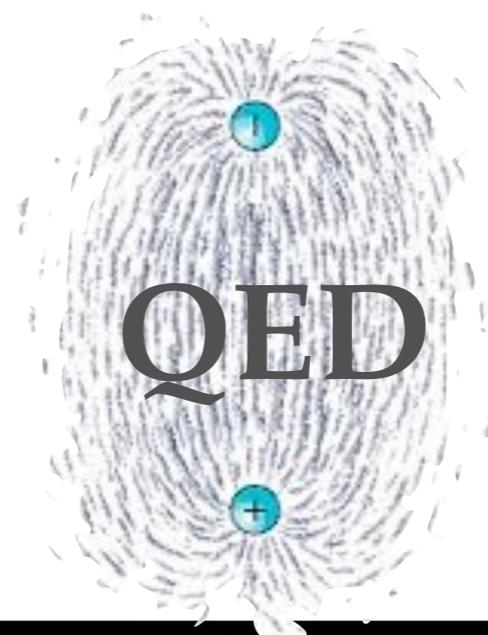
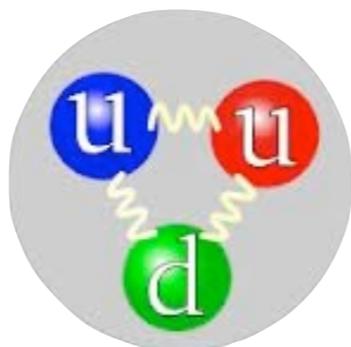
$\lambda = \frac{\pi}{k}$

# Conformal Bootstrap and QFT

QFTs are on RG flow  
between (or nearly  
between) CFTs

E.g. The Standard Model  
QCD

???



$$m_Z^{-1} \sim 10^{-18} m$$

$$\Lambda_{\text{QCD}}^{-1} \sim 10^{-15} m$$

$$m_e^{-1} \sim 10^{-12} m$$